Regulatory mechanism of length-dependent activation in skinned porcine ventricular muscle: role of thin filament cooperative activation in the Frank-Starling relation

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Cardiac sarcomeres produce greater active force in response to stretch, forming the basis of the Frank-Starling mechanism of the heart. The purpose of this study was to provide the systematic understanding of length-dependent activation by investigating experimentally and mathematically how the thin filament “on–off” switching mechanism is involved in its regulation. Porcine left ventricular muscles were skinned, and force measurements were performed at short (1.9 µm) and long (2.3 µm) sarcomere lengths. We found that 3 mM MgADP increased Ca²⁺ sensitivity of force and the rate of rise of active force, consistent with the increase in thin filament cooperative activation. MgADP attenuated length-dependent activation with and without thin filament reconstitution with the fast skeletal troponin complex (sTn). Conversely, 20 mM of inorganic phosphate (Pi) decreased Ca²⁺ sensitivity of force and the rate of rise of active force, consistent with the decrease in thin filament cooperative activation. Pi enhanced length-dependent activation with and without sTn reconstitution. Linear regression analysis revealed that the magnitude of length-dependent activation was inversely correlated with the rate of rise of active force. These results were quantitatively simulated by a model that incorporates the Ca²⁺-dependent on–off switching of the thin filament state and interfilament lattice spacing modulation. Our model analysis revealed that the cooperativity of the thin filament on–off switching, but not the Ca²⁺-binding ability, determines the magnitude of the Frank-Starling effect. These findings demonstrate that the Frank-Starling relation is strongly influenced by thin filament cooperative activation.

INTRODUCTION

At the turn of the 20th century, Frank and Starling discovered that cardiac pump function is enhanced as ventricular filling is increased (i.e., the Frank-Starling law of the heart; see Katz, 2002 and references therein). The “law” forms the fundamental principle of the heart in cardiovascular physiology, defining the relation between the diastolic and systolic performances of cardiac chambers. It is widely accepted that the length dependence of Ca²⁺-based myofibrillar activation (i.e., expressed as “Ca²⁺ sensitivity of force”) largely underlies the law (e.g., Allen and Kurihara, 1982; Allen and Kentish, 1985; Kentish et al., 1986); however, the molecular mechanism of this seemingly simple phenomenon still remains elusive and warrants an in-depth investigation.

The cross-bridge formation is a stochastic process in the striated muscle sarcomere (e.g., Huxley, 1957). Therefore, it has been proposed that the binding of myosin to actin is enhanced upon the reduction in the distance between the thick and thin filaments (i.e., interfilament lattice spacing), resulting in an increase in active force production and, apparently, Ca²⁺ sensitivity of force (Ishiwata and Oosawa, 1974; McDonald and Moss, 1995; Fuchs and Wang, 1996; Fukuda et al., 2000). Indeed, studies with synchrotron x ray revealed that passive force due to extension of the giant elastic protein titin (also known as connectin) modulates the lattice spacing within the physiological sarcomere length (SL) range in cardiac muscle (Cazorla et al., 2001; Fukuda et al., 2003, 2005). Konhilas et al. (2002b), however, challenged this proposal, demonstrating that the lattice spacing and Ca²⁺ sensitivity of force are not well correlated. Therefore, factors other than the titin-based lattice spacing modulation are likely at play in the regulation of length-dependent activation.

As has been reported, multiple cooperative processes are involved in active force generation in striated muscle...
(e.g., Brandt et al., 1982, 1987, 1990; Moss et al., 1985); i.e., cooperative binding of Ca^{2+} to troponin (Tn) C (TnC), cooperative binding of myosin to the thin filaments, and synergistic interactions between myosin binding to actin and Ca^{2+} binding to TnC (e.g., Bremel et al., 1973; Güth and Potter, 1987; Hoar et al., 1987; Zot and Potter, 1989; Swartz and Moss, 1992). Likewise, it is widely accepted that the formation of strongly bound cross-bridges enhances cooperative recruitment of neighboring myosin to the thin filaments. Bremel and Weber (1972) were the first to demonstrate in solution that an increase in the fraction of rigor cross-bridges (rigor myosin subfragment 1) cooperatively activates myosin ATPase, as with the increased Ca^{2+} concentration, indicating that Ca^{2+} and strongly bound cross-bridges synergistically regulate the “on–off” equilibrium of the thin filament state. Later, the group of Moss provided evidence in skinned muscle fibers that actomyosin interaction is indeed promoted in the presence of the strong-binding cross-bridge analogue N-ethylmaleimide myosin subfragment 1 (NEM-S1) via activation of the thin filaments, as manifested by the increased rate of contraction (Swartz and Moss, 1992, 2001; Fitzsimons et al., 2001a,b) and the increased Ca^{2+} sensitivity of force (Fitzsimons and Moss, 1998; Fitzsimons et al., 2001a,b). Similarly, we previously reported that the application of MgADP, i.e., the ensuing formation of the actomyosin–ADP complex, cooperatively enhances cross-bridge recruitment and that of inorganic phosphate (Pi) does the opposite in skinned fibers of cardiac and skeletal muscles (Shimizu et al., 1992; Fukuda et al., 1998, 2000, 2001a).

Earlier studies have suggested that the thin filament–based on–off switching mechanism is involved in the regulation of length-dependent activation in cardiac muscle. Indeed, Fitzsimons and Moss (1998) reported that length-dependent activation is attenuated in the presence of NEM-S1. Fukuda et al. (2000) confirmed this notion by providing evidence that the length dependence becomes smaller in the presence of MgADP (hence the actomyosin–ADP complex). It has also been reported that direct modulation of thin filament regulatory proteins, e.g., Tn isoform changes (Arteaga et al., 2000; Terui et al., 2008) or TnI point mutation (Tachampa et al., 2007), markedly affects length-dependent activation. Therefore, it is likely that length-dependent activation depends on the state of the thin filaments, either modulated directly by regulatory protein isoform switching or indirectly by strongly bound cross-bridges. However, it is still unknown how Ca^{2+}-dependent on–off switching of the thin filament state and interfilament lattice spacing coordinate to regulate myocardial length-dependent activation.

Accordingly, this study was undertaken to systematically uncover the molecular basis of length-dependent activation in cardiac muscle, focusing on the role of thin filament cooperative activation. We varied the level of thin filament cooperative activation in skinned porcine left ventricular muscle (PLV) directly by Tn exchange or indirectly by the application of MgADP or Pi, or the combination of both. Our analysis revealed that the magnitude of length-dependent activation is inversely related to the rate of rise of active force, highlighting a pivotal role of thin filament cooperative activation in the regulation of the Frank-Starling relation. Furthermore, our mathematical model analyses revealed the relationship between the characteristics of thin filament activation and length-dependent activation, and led us to conclude that length-dependent activation is under the strong control of thin filament cooperative activation.

**MATERIALS AND METHODS**

All experiments performed in this study conform to the Guide for the Care and Use of Laboratory Animals (1996. National Academy of Sciences, Washington D.C.). For expanded materials and methods, please see the supplemental material.

**Preparation of skinned muscle**

Skinned muscles were prepared according to the method in our recent studies (Terui et al., 2008; Matsuba et al., 2009). In brief, porcine hearts (from ~1.0-yr-old animals) were obtained from a local slaughterhouse. Muscle strips (1–2 mm in diameter and ~10 mm in length) were dissected from the papillary muscle of the left ventricle in Ca^{2+}-free Tyrode’s solution (see Fukuda et al., 2000a for composition) containing 30 mM 2,3-butanediol monoxime (BDM).

**Muscle mechanics**

Isometric force was measured according to the method in our recent studies (Terui et al., 2008; Udaka et al., 2008; Matsuba et al., 2009). In brief, PLVs were skinned in relaxing solution (5 mM MgATP, 40 mM BES, 1 mM Mg^{2+}, 10 mM EGTA, 1 mM dithiothreitol, 15 mM phosphocreatine, 15 U/ml creatine phosphokinase (CPK), and 180 mM ionic strength [adjusted by K-propionate], pH 7.0), containing 1% (wt/vol) Triton X-100 and 10 mM BDM overnight at ~3°C (Fukuda et al., 2003, 2005). Muscles were stored for up to 3 wk at ~2°C in relaxing solution containing 50% (vol/vol) glycerol. All solutions contained protease inhibitors (0.5 mM PMSF, 0.04 mM leupeptin, and 0.01 mM EDTA).

Small thin preparations (~100 µm in diameter and ~2 mm in length) were dissected from the porcine ventricular strips for force measurement. SL was measured by laser diffraction during relaxation, and active and passive forces were measured at 15°C (pCa adjusted by Ca^{2+}/EGTA based on a computer program by Fabiato, 1988). MgADP (up to 10 mM) or Pi (up to 20 mM) was added to the individual activating solutions in accordance with our previous studies (Fukuda et al., 2000, 2001a), while maintaining ionic strength at 180 mM. When MgADP was present, 0.1 mM P_i P_i di(adenosine-5’)-pentaphosphate was added to both activating and relaxing solutions, with no CPK to maintain the ADP/ATP ratio (Fukuda et al., 1998, 2000). We also used pimobendan (provided by Nippon Boehringer Ingelheim) to increase the affinity of TnC for Ca^{2+} (Fukuda et al., 2000). Pimobendan was initially dissolved in DMSO and diluted with the individual solutions (Fukuda et al., 2000). The final concentration of DMSO was 1%, having no effect on active or passive force, as observed in our previous study (Fukuda et al., 2000).
The muscle preparation was first immersed in relaxing solution, and SL was set at 1.9 µm. Active and passive forces were measured at 1.9 µm and then at 2.3 µm, as described in our previous studies (rundown <10% for active and passive forces; Fukuda et al., 2003, 2005; Terui et al., 2008). Active force data were fitted to the Hill equation (Fukuda et al., 2000), and the difference between the values of the midpoint of the force-pCa curve (i.e., pCa50) at SL 1.9 and 2.3 µm was used as an index of the SL dependence of Ca2+ sensitivity of force (expressed as ΔpCa50). The steepness of the force-pCa curve was expressed as the Hill coefficient (nH).

The rate of rise of active force was assessed according to the method in our previous work (Fukuda et al., 2001b), using the preparations used for the steady-state isometric force measurement. In brief, SL was set at 1.9 µm in relaxing solution. The preparation was then immersed in low EGTA (0.5 mM) relaxing solution for 1 min and transferred to the control activating solution. In brief, SL was set at 1.9 µm in relaxing solution. The muscle preparation was first immersed in relaxing solution. The preparation was then immersed in low EGTA (0.5 mM) relaxing solution for 1 min and transferred to the control activating solution. In brief, SL was set at 1.9 µm in relaxing solution. The preparation was then immersed in low EGTA (0.5 mM) relaxing solution for 1 min and transferred to the control activating solution.

Model analysis

The model calculates the active isometric force at a given SL and at a given Ca2+ concentration, based on the SL-dependent change in the lattice spacing and the Ca2+-based on-off switching of the thin filament state. The on-off state was defined according to the lateral fluctuation of the thin filaments in the myofilament lattice (Ishiwata and Fujime, 1972; Umezume and Fujime, 1975; Yoshino et al., 1978; Yanagida et al., 1984; see Fig. S1); however, the lateral fluctuation does not necessarily represent the physical thermal fluctuation of the thin filaments, but rather, it mathematically portrays the equilibrium of the thin filament state between “off” and “on,” depending on the Ca2+ concentration (Solaro and Rarick, 1998).

In our model, the overlap length between the thick and thin filaments at a SL of L is given by:

\[
\frac{1}{2}(L_0 - L),
\]

where \(L_0\) (3.8 µm) is the maximal SL at no filament overlap. Thick and thin filament length was assumed to be 1.6 and 1.1 µm, respectively (Sosa et al., 1994). In this study, the overlap length was set to be constant (0.75 µm), independent of SL because in the SL range between 1.9 and 2.3 µm (where the experiments were performed), the number of myosin heads in the overlap region reportedly remains constant based on the thick filament geometry (Sosa et al., 1994).

Next, we assumed that the lattice spacing, \(d\), decreases upon the increase in SL under the constant lattice volume, \(V\), as has been observed in x-ray diffraction studies (Cazorla et al., 2001; Fukuda et al., 2003, 2005), according to the following equation:

\[
V = d^2 \cdot L.
\]

Based on the value of \(d_0\) (i.e., the \(d_0\) lattice spacing) of 43 nm at SL 2.0 µm (see Fukuda et al., 2003 for the \(d_0\) value of sarco-meres expressing both N2B and N2BA titins at similar levels, as in PLV; Terui et al., 2008), the lattice spacing is estimated to be 28.7 nm, and thereby the lattice volume, \(V\), was set to be 0.0016 µm3.

Next, we described the position-dependent probability of actomyosin interaction by the Gaussian distribution (Ishiwata and Oosawa, 1974):

\[
P(q) = \frac{1}{\sqrt{\pi \sigma_q^2}} \exp\left(-\frac{q^2}{\sigma_q^2}\right),
\]

where \(q\) is a lateral coordinate perpendicular to the filament long axis, and \(\sigma_q\) is the width of the Gaussian distribution (variance, \(\sigma_q^2/2\)). To take into account the Ca2+-dependent change of the actomyosin interaction, the degree of \(\sigma_q\) was changed in accordance with the Ca2+ concentration based on the Hill equation (see Ishiwata and Oosawa, 1974 and Figs. S1 and S2):

\[
\sigma_q(pCa) = \sigma_{\text{max}} \frac{1}{1 + 10^{\frac{pCa-pCa_{50}}{\phi_{\text{H_actin}}(pCa_{50}-pCa_{50})}}},
\]

where \(\sigma_{\text{max}}\) is the maximal width of the Gaussian distribution (21 nm), determining the maximal interaction probability at the saturating Ca2+ concentration (pCa 4.5). The parameter \(\phi_{\text{H_actin}}\) represents the cooperativity of thin filament activation in the model calculation, and \(pCa_{50}\) represents the sensitivity of the thin filaments to Ca2+.

Here, we considered that the actomyosin interaction takes place over the region where the lateral coordinate, \(q\), exceeds a certain distance, \(d-a\), because myosin heads are located apart from the thick filament backbone up to a distance \(a\) (24 nm). The cumulative interaction probability with respect to the unit overlap length, \(I\), is therefore given by:

\[
I = \int_{d-a} \frac{1}{d-a} P(q) dq.
\]

Finally, the active isometric force at a given SL and pCa is expressed as the product of overlap length and interaction probability:

\[
F = F_0 \cdot \frac{1}{2}(L_0 - L) \cdot \frac{1}{\sqrt{\pi}} \int_{d-a} \frac{1}{\sigma_q(pCa)} \exp\left(-\frac{q^2}{\sigma_q^2(pCa)}\right) dq,
\]
where $F_0$ (45 and 36, for control and sTn-reconstituted PLV, respectively) is the fitting parameter to quantitatively simulate the experimental results.

**Statistics**
Significant differences were assigned using the paired or unpaired Student’s t test as appropriate. Data are expressed as mean ± SEM, with $n$ representing the number of muscles. Linear regression analyses were performed in accordance with the method described in previous studies (Fukuda et al., 2001b; Terui et al., 2008). Statistical significance was assumed to be P < 0.05. NS indicates $P > 0.05$.

**Online supplemental material**
The supplemental material provides an expanded description of our model analysis. In addition, Fig. S1 shows a schematic illustration of our model used to simulate the present experimental data. Fig. S2 provides characteristics of our model, showing how Ca$^{2+}$ sensitivity of force is changed in response to a change in thin filament cooperative activation. Fig. S3 shows the relation of Ca$^{2+}$ binding to the thin filaments or thin filament cooperative activation versus length-dependent activation in our model. Fig. S4 shows the experimentally observed effect of MgADP on the rate of active force redevelopment, $k_c$ (overall cross-bridge cycling rate; see Discussion), at varying activation levels. Fig. S5 shows the experimentally obtained relation between SL and active force at various Ca$^{2+}$ concentrations (converted from force–pCa curves) and the simulation by our model. Fig. S6 shows the model simulation showing the effect of an increase in the average length of myosin heads on Ca$^{2+}$ sensitivity of force and length-dependent activation. The online supplemental material is available at http://www.jgp.org/cgi/content/full/jgp.201010502/DC1.

**RESULTS**

**Effect of MgADP or Pi on the rate of rise of active force**
First, we investigated the effect of various concentrations of MgADP or Pi on the rate of rise of active force under the control condition without sTn reconstitution. We found that MgADP significantly decreased $t^{1/2}$ at low concentrations (1 and 3 mM) but increased it at a high concentration (10 mM) (Fig. 1 A), and Pi exerted apparently similar effects at low (1, 3, and 5 mM) and high (10 and 20 mM) concentrations (Fig. 1 B). The accelerating effect of MgADP was maximal at 3 mM, whereas the decelerating effect of Pi reached a quasi-plateau at 10 mM. Therefore, based on the previous studies indicating that the rate of contraction is modulated by a change in the fraction of strongly bound cross-bridges via alteration of the on–off

Figure 1. Effect of MgADP or Pi on the rate of rise of active force. pCa 4.5; SL, 1.9 µm. (A) Effect of MgADP. (Left) Typical chart tracings showing active force responses in the absence and presence of 3 and 10 mM MgADP (in the same preparation). (Right) Graph summarizing the effects of various concentrations of MgADP on $t^{1/2}$. *, $P < 0.05$ compared with 0 mM MgADP. Active force compared with 0 mM MgADP (in percent) was 99.87 ± 3.74, 97.37 ± 5.57, 93.21 ± 5.44, and 90.07 ± 3.27, with 1, 3, 5, and 10 mM MgADP, respectively; $n = 5$. (B) Effect of Pi. (Left) Typical chart tracings showing active force responses in the absence and presence of 3 and 20 mM Pi (in the same preparation). (Right) Graph summarizing the effects of various concentrations of Pi on $t^{1/2}$. *, $P < 0.05$ compared with 0 mM Pi. Active force compared with 0 mM MgADP (in percent) was 99.87 ± 3.74, 97.37 ± 5.57, 93.21 ± 5.44, and 90.07 ± 3.27, with 1, 3, 5, and 10 mM MgADP, respectively; $n = 5$. In both A and B, arrowheads and double arrowheads indicate the points at which solution was switched from low EGTA (0.5 mM) relaxation to contraction and from contraction to high EGTA (10 mM) relaxation, respectively. The time to half-maximal activation (50%) was measured as indicated by $\alpha$ and $\beta$, and the relative value, i.e., $\beta/\alpha$, was obtained for each preparation and defined as $t^{1/2}$. Note that $t^{1/2}$ is $\sim$1.0 in the absence of MgADP or Pi, indicating reproducibility of the rate of rise of active force. (C) Relation between $t^{1/2}$ and $V^{1/2}$ obtained in the presence of varying concentrations of MgADP (left, 0–10 mM) and Pi (right, 0–20 mM). A significant linear relationship existed for both MgADP ($R = 0.68$; $P < 0.0005$) and Pi ($R = 0.70$; $P < 0.0001$). Data taken from bar graphs in A and B.
equilibrium of the thin filament state (Swartz and Moss, 1992, 2001; Fitzsimons et al., 2001a,b), we regarded 3 mM MgADP as the amount to enhance thin filament cooperative activation and 20 mM Pi to reduce it and used them in the following experiments.

Fig. 1 C shows the relationship between $t_{1/2}$ and $V_{1/2}$. We found that a significant linear relationship with a similar slope value existed between the parameters in the presence of varying concentrations of MgADP (Fig. 1 C, left; slope 1.07) or Pi (right, slope 1.24), indicating that $t_{1/2}$ reflects the rate of rise of active force.

Effect of MgADP or Pi on length-dependent activation with and without sTn reconstitution

Next, we investigated how the SL-dependent increase in Ca\(^{2+}\) sensitivity of force responds to alteration of thin filament cooperative activation. In this series of experiments, we performed sTn reconstitution to directly enhance thin filament cooperative activation, as demonstrated in our previous study (Terui et al., 2008), with and without MgADP or Pi.

We found that 3 mM MgADP or sTn reconstitution similarly shifted the force–pCa curve leftward, to a greater magnitude at SL 1.9 µm than at 2.3 µm, and consequently decreased $\Delta pC_{a_{50}}$ (Fig. 2, A and B). Ca\(^{2+}\) sensitivity of force was synergistically increased by MgADP in sTn-reconstituted PLV, accompanied by a marked attenuation of length-dependent activation (Fig. 2 B). As shown in Fig. 2 C, the rate of rise of active force was increased by MgADP or sTn reconstitution by a similar magnitude. Similar to the finding on Ca\(^{2+}\) sensitivity of force, the rate of rise of active force was increased by MgADP in sTn-reconstituted PLV (Fig. 2 C).

We then tested the effect of Pi on length-dependent activation. Without sTn reconstitution, 20 mM Pi shifted the force–pCa curve rightward to a greater magnitude at SL 1.9 µm than at 2.3 µm, and consequently increased $\Delta pC_{a_{50}}$ (Fig. 3 A). Pi increased $\Delta pC_{a_{50}}$ also in sTn-reconstituted PLV (Fig. 3 B). In contrast to MgADP, Pi retarded the rate of rise of active force in both control and sTn-reconstituted PLV (Fig. 3 C). The values of $pC_{a_{50}}$, $nH$, and maximal force at SL 1.9 and 2.3 µm under various conditions are summarized in Table I.

Fig. 4 summarizes the relationship between $t_{1/2}$, $pC_{a_{50}}$ and $\Delta pC_{a_{50}}$ obtained with MgADP or Pi in control and sTn-reconstituted PLV. We found that $pC_{a_{50}}$ and $\Delta pC_{a_{50}}$ were linearly correlated with each other (Fig. 4 A), and that $pC_{a_{50}}$ and $\Delta pC_{a_{50}}$ were a linear function of $t_{1/2}$ (Fig. 4, B and C). As reported previously (Dobesh et al., 2002), however, no significant correlation was found between $nH$ and $pC_{a_{50}}$ or $\Delta pC_{a_{50}}$ (Fig. 5).

Effect of pimobendan on length-dependent activation with and without sTn reconstitution

The observed relationship of Ca\(^{2+}\) sensitivity of force and length-dependent activation may be a consequence...
associated with the leftward shift of the force–pCa curve (Hanft et al., 2008). We therefore tested the effect of pimobendan on length-dependent activation, with passive force carefully controlled (which was not performed in our previous study on rat ventricular trabeculae; Fukuda et al., 2000), because the compound has been reported to specifically increase the affinity for Ca\(^{2+}\) of the low-affinity site of TnC (see Hagemeijer, 1993 and references therein).

Pimobendan (2 × 10\(^{-4}\) M) shifted the force–pCa curve leftward to a magnitude similar to that by MgADP or sTn reconstitution at SL 1.9 μm (i.e., ~0.2 pCa units; see Fig. 2), with no effect on passive force (Fig. 6 A). However, unlike MgADP or sTn reconstitution, pimobendan exerted no effect on ΔpCa\(_{50}\) (see Fukuda et al., 2000). The Ca\(^{2+}\)-sensitizing effect of pimobendan was markedly diminished after sTn reconstitution, with no significant increase in Ca\(^{2+}\) sensitivity of force at either SL (Fig. 6 B). Likewise, pimobendan did not affect the rate of rise of active force in control or sTn-reconstituted PLV (Fig. 6 C).

The values of pCa\(_{50}\), n\(_{H}\), and maximal force obtained with and without pimobendan are summarized in Table II.

**Simulation of length-dependent activation**

Finally, we analyzed the experimental findings based on a mathematical model (Ishiwata and Oosawa, 1974; Shimamoto et al., 2007; refer to Materials and methods and supplemental material for details). In this model, active isometric force is given by the interaction probability between the thick and thin filaments, and the probability depends on two factors: (1) Ca\(^{2+}\) concentration and (2) interfilament lattice spacing (Fig. S1). The equilibrium of the thin filament state between “off” and “on” (see Solaro and Rarick, 1998 and references therein) was assumed to change with the Ca\(^{2+}\) concentration based on the Hill equation, and expressed as lateral fluctuation in the myofilament lattice (Figs. S1 and S2). We performed experiments within the SL range (i.e., 1.9–2.3 μm); the overlap length between the thick and thin filaments is considered not to change significantly (see Moss and Fitzsimons, 2002 and references therein), but the lattice spacing does, due to titin extension, as revealed by previous studies with muscles expressing both N2B and N2BA titins (Fukuda et al., 2003), as in PLV (Terui et al., 2008).

Fig. 7 A shows the force–pCa curves simulated by our model for the experimental data with and without sTn reconstitution in PLV. The model parameters are n\(_{H, \text{actin}}\) and pCa\(_{50, \text{actin}}\), representing the characteristics of thin filament on–off switching in response to Ca\(^{2+}\). Based on Eq. 6, we simulated the force–pCa curves of PLV with and without sTn reconstitution, and thereby the pCa\(_{50}\) and ΔpCa\(_{50}\) values were calculated (for optimization of fitting, see supplemental material).

Under the control condition without sTn reconstitution, a reduction in the lattice spacing due to an increase in SL from 1.9 to 2.3 μm (refer to Materials and methods) increased maximal Ca\(^{2+}\)-activated force and shifted the force–pCa curve leftward, resulting in ΔpCa\(_{50}\) of 0.24 pCa units. As shown in Fig. S2, an increase in n\(_{H, \text{actin}}\)
decreased \Delta pC_{50} and concomitantly shifted the midpoint of the force–pCa curve rightward. On the other hand, an increase in \( pC_{50,\text{actin}} \) linearly shifted the force–pCa curve leftward. To reproduce the experimental data after sTn reconstitution, we increased both \( n_{H,\text{actin}} \) and \( pC_{50,\text{actin}} \) as a result, the attenuation of length-dependent activation was well simulated, accompanied by appropriate \( pC_{50} \) values for both SLs (as in Fig. 7 A; compare Figs. 2 and 3, and Terui et al., 2008).

Finally, we systematically investigated how varying the values of \( n_{H,\text{actin}} \) and \( pC_{50,\text{actin}} \) affects length-dependent activation by constructing a 3-D graph consisting of \( n_{H,\text{actin}} \) and \( pC_{50,\text{actin}} \) (Fig. 7 B). We found that \( \Delta pC_{50} \) was strongly influenced by \( n_{H,\text{actin}} \). However, the contribution of \( pC_{50,\text{actin}} \) to \( \Delta pC_{50} \) was minimal throughout the range we examined (see also Fig. S3). We plotted the pairs of the values of \( n_{H,\text{actin}} \) and \( pC_{50,\text{actin}} \) that fulfill the linear relationship between \( pC_{50} \) and \( \Delta pC_{50} \) (i.e., red points in Fig. 7 B), which was experimentally obtained in Fig. 4 A.

### DISCUSSION

We demonstrated in this study that the Frank-Starling relation is strongly influenced by thin filament cooperative activation. The SL-dependent increase in \( Ca^{2+} \) sensitivity of force was inversely related to the rate of rise of active force, suggesting that length-dependent activation is tuned via on-off switching of the thin filament state in cardiac muscle. Further, our model analysis revealed that thin filament cooperative activation, but not the affinity for \( Ca^{2+} \), determines the magnitude of length-dependent activation, coupled with lattice spacing modulation. Here, we discuss the present findings, focusing on the role of thin filament cooperative activation in the regulation of length-dependent activation in cardiac muscle.

As is well established, strongly bound cross-bridges cooperatively activate the thin filaments (e.g., Bremel and Weber, 1972), resulting in the promotion of actomyosin interaction (refer to Introduction). Previous studies showing that NEM-S1 increases the speed of contraction in skinned fibers (Swartz and Moss, 1992, 2001; Fitzsimons et al., 2001a,b) support the notion that strongly bound cross-bridges accelerate the recruitment of neighboring myosin to the thin filaments via enhanced thin filament cooperative activation. In the present study, MgADP decreased \( t_{1/2} \) at low concentrations (1 and 3 mM) but increased it at a high concentration (10 mM), and \( Pi \) exerted apparently similar effects at low (1, 3, and 5 mM) and high (10 and 20 mM) concentrations (Fig. 1; see Kentish, 1986). However, the underlying molecular mechanisms for the modulation of \( t_{1/2} \), i.e., the rate of rise of active force, should differ for MgADP and \( Pi \). At low MgADP concentrations (e.g., 3 mM used in the present study), the actomyosin–ADP complex may exhibit its promoting effect on actomyosin interaction by enhancing thin filament cooperative activation, but at high MgADP concentrations, the presence of large fractions of the complex may cause deceleration of contraction due to its slow cycling rate, as demonstrated

### TABLE I

Summary of the values of passive force, maximal active force, \( pC_{50} \), and \( n_{H} \) in PLV under various conditions

<table>
<thead>
<tr>
<th>SL (µm)</th>
<th>Passive force (mN/mm²)</th>
<th>Maximal force (mN/mm²)</th>
<th>( pC_{50} )</th>
<th>( \Delta pC_{50} )</th>
<th>( n_{H} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>Without sTn reconstitution</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Control 1.9</td>
<td>57.22 ± 2.29</td>
<td>5.56 ± 0.01</td>
<td>3.86 ± 0.13</td>
<td></td>
<td></td>
</tr>
<tr>
<td>2.3</td>
<td>75.79 ± 5.72</td>
<td>5.80 ± 0.01</td>
<td>3.17 ± 0.24</td>
<td></td>
<td></td>
</tr>
<tr>
<td>+ADP 1.9</td>
<td>67.19 ± 1.39* (78.41 ± 2.97)</td>
<td>5.82 ± 0.01*</td>
<td>3.44 ± 0.10*</td>
<td></td>
<td></td>
</tr>
<tr>
<td>2.3</td>
<td>75.61 ± 2.64</td>
<td>5.93 ± 0.01*</td>
<td>3.06 ± 0.16</td>
<td></td>
<td></td>
</tr>
<tr>
<td>+Pi 1.9</td>
<td>21.66 ± 1.39* (44.58 ± 5.53)</td>
<td>5.29 ± 0.03*</td>
<td>3.74 ± 0.31</td>
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<td>2.3</td>
<td>38.53 ± 2.48*</td>
<td>5.60 ± 0.03*</td>
<td>4.18 ± 0.33*</td>
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<td>With sTn reconstitution</td>
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<td>Control 1.9</td>
<td>51.83 ± 1.84</td>
<td>5.79 ± 0.01*</td>
<td>3.20 ± 0.18*</td>
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<td>2.3</td>
<td>58.76 ± 5.14*</td>
<td>5.93 ± 0.02*</td>
<td>3.00 ± 0.08</td>
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<tr>
<td>+ADP 1.9</td>
<td>44.75 ± 5.61 (48.07 ± 4.49)</td>
<td>5.94 ± 0.02*</td>
<td>3.38 ± 0.18</td>
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<tr>
<td>2.3</td>
<td>49.90 ± 4.76*</td>
<td>6.01 ± 0.02*</td>
<td>3.52 ± 0.19</td>
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<tr>
<td>+Pi 1.9</td>
<td>22.36 ± 1.19* (57.17 ± 2.70)</td>
<td>5.50 ± 0.04*</td>
<td>2.39 ± 0.29*</td>
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<tr>
<td>2.3</td>
<td>44.76 ± 2.32*</td>
<td>5.74 ± 0.04*</td>
<td>2.93 ± 0.30</td>
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Data are for Figs. 2 and 3. Maximal force was obtained by activating muscle at pCa 4.5 before construction of the force–pCa curve at each SL (passive force was measured just before activation at pCa 4.5). Numbers in parentheses indicate maximal force values obtained before ADP or Pi application. Maximal force obtained before sTn reconstitution at SL 1.9 µm: 54.36 ± 2.80, 48.65 ± 3.66, and 61.35 ± 2.08 mN/mm² for the control, ADP-, and Pi-treated group, respectively (\( P > 0.05 \) compared with the value after sTn reconstitution).

\*\( P < 0.05 \) compared with the prior value.

\*\( P < 0.05 \) compared with the corresponding values in control group with sTn reconstitution.

\*\( P < 0.05 \) compared with the corresponding values in control group with sTn reconstitution.
previously in experiments measuring the shortening of the velocity at zero load (e.g., Cooke and Pate, 1985; Metzger, 1996), despite the highly activated state of the thin filaments. On the other hand, the binding of Pi to the actomyosin complex is reportedly enhanced upon the increase in the strain of the complex (Webb et al., 1986; Metzger, 1996). Therefore, at low concentrations, Pi may preferentially bind to the slowly cycling actomyosin–ADP complex, resulting in an increase in the rate of rise of active force, as demonstrated previously in experiments measuring kinetics following flash photolysis (Lu et al., 1993; Araujo and Walker, 1996), the shortening of the velocity at zero load (Metzger, 1996) and the rate of force redevelopment ($k_{tr}$) (Tesi et al., 2000). However, at high concentrations (e.g., 20 mM used in the present study), Pi may decrease the fraction of the actomyosin–ADP complex to a level where neighboring myosin cannot be effectively recruited to actin, resulting in a decrease in the rate of rise of active force. Therefore, although the alteration of $t_{1/2}$ includes processes other than thin filament cooperative activation, the present findings allow us to consider that it at least in part reflects a change in thin filament cooperative activation.

MgADP at 3 mM increased the rate of rise of active force and, concomitantly, left-shifted the force–pCa curve (Figs. 1 and 2). The magnitude of the change was similar to that observed upon sTn reconstitution (i.e., direct modulation of regulatory proteins to enhance thin filament cooperative activation; see Terui et al., 2008) for both $t_{1/2}$ and Ca$^{2+}$ sensitivity of force (Fig. 2). These findings suggest that, albeit modulated via different pathways, i.e., either indirectly or directly, thin filament cooperative activation is enhanced by a similar magnitude with 3 mM MgADP and sTn reconstitution. Interestingly, 3 mM MgADP increased both Ca$^{2+}$ sensitivity of force and the rate of rise of active force in sTn-reconstituted PLV, accompanied by a marked depression of length-dependent activation (Fig. 2). These additive effects of MgADP suggest that thin filament cooperative activation can be synergistically modulated via strong-binding cross-bridge formation and regulatory protein isoform switching. In contrast, 20 mM Pi exerted effects opposite to those of 3 mM MgADP, with and without sTn reconstitution, by decreasing Ca$^{2+}$ sensitivity of force and slowing the rate of rise of active force (Fig. 3). Therefore,

Figure 4. Linear regression analyses between $t_{1/2}$, pCa$_{50}$, and $\Delta$pCa$_{50}$. The following relationships are shown: pCa$_{50}$ versus $\Delta$pCa$_{50}$ (A, $R = 0.98$; $P < 0.0005$), $t_{1/2}$ versus pCa$_{50}$ (B, $R = 0.96$; $P < 0.005$), and $t_{1/2}$ versus $\Delta$pCa$_{50}$ (C, $R = 0.97$; $P < 0.001$). Plots were constructed using the data in Figs. 2 and 3 (pCa$_{50}$ obtained at SL 1.9 µm).

Figure 5. Relation between the $n_{H}$ of the force–pCa curve (i.e., $n_{H}$ [force]) and $\Delta$pCa$_{50}$ obtained experimentally under various conditions. (A) Plot of $n_{H}$ (force) versus pCa$_{50}$. (B) Plot of $n_{H}$ (force) versus $\Delta$pCa$_{50}$. No significant correlation was found between the parameters in either graph (see Dobesh et al., 2002). For both A and B, the $n_{H}$ values used were from the force–pCa curve at SL 1.9 µm in Figs. 2 and 3 (see Table I). NS, not significant.
the observed effect of 20 mM Pi on length-dependent activation likely results from the reduced thin filament cooperative activation. Furthermore, pimobendan did not affect length-dependent activation, indicating that Ca$^{2+}$ binding to TnC is not the parameter determining the magnitude of this phenomenon. Therefore, given the close relationship between $t_{1/2}$ (or pCa$^{50}$) and $\Delta$pCa$^{50}$ (Fig. 4), we consider that thin filament cooperative activation plays a pivotal role in setting the magnitude of length-dependent activation.

It has been reported in various experimental settings that a positive feedback mechanism exists between Ca$^{2+}$ binding to TnC and cross-bridge formation in the sarcosome (Allen and Kurihara, 1982; Güth and Potter, 1987; Hoar et al., 1987; Hofmann and Fuchs, 1988; Zot and Potter, 1989). Therefore, the linear relationship of $t_{1/2}$ versus pCa$^{50}$ observed in the present study (Fig. 4) likely reflects the positive feedback effect on Ca$^{2+}$ binding to TnC via cross-bridge formation due to enhanced thin filament cooperative activation.

One may point out that $k_{tr}$, i.e., the sum of the apparent rate of cross-bridge attachment ($f_{app}$) and detachment ($g_{app}$) (Brenner, 1988; Swartz and Moss, 1992; Fitzsimons et al., 2001a,b, and references therein; Terui et al., 2008), more suitably represents thin filament cooperative activation in muscle mechanics than $t_{1/2}$. Indeed, $k_{tr}$ is reportedly increased upon enhanced thin filament cooperative activation, when modulated directly (e.g., sTn reconstitution; Terui et al., 2008) or indirectly (NEM-S1 application; Swartz and Moss, 1992; Fitzsimons et al., 2001a,b). In the present study, 3 mM MgADP decreased $k_{tr}$ at both maximal and submaximal activations (Fig. S4), in agreement with the result of previous studies with rabbit skeletal muscle (Lu et al., 1993; Tesi et al., 2000). MgADP is known to inhibit the release of ADP from the actomyosin complex at the end of the cross-bridge cycle (see Fukuda et al., 1998, 2000, and references therein). Therefore, the inhibitory effect of MgADP on $g_{app}$ may overshadow its accelerating effect on cross-bridge formation ($f_{app}$), resulting in a decrease in $k_{tr}$. However, the observed increase in the rate of rise of active force, regardless of the Tn isoform, suggests that MgADP at low concentrations (such as 3 mM in the present experimental setting; see Figs. 1 and 2) accelerates $f_{app}$ via enhancement of thin filament cooperative activation.

Pimobendan did not significantly increase Ca$^{2+}$ sensitivity of force after sTn reconstitution in PLV (Fig. 6). It has, however, been reported that pimobendan increases Ca$^{2+}$ sensitivity of force in amphibian skeletal muscles (Piazzesi et al., 1987; Wakisaka et al., 2000). The absence of Ca$^{2+}$ sensitization observed in the present study may indicate the compound’s specificity regarding the site of action; namely, in mammals, the binding affinity of pimobendan for (fast) skeletal TnC may be lower than that for cardiac TnC, producing a minimal effect on (fast) skeletal muscle. It should also be pointed out that
However, it is unclear to what extent the steady-state nH high Ca\textsuperscript{2+} concentrations, MgADP application, or sTn reconstitution results not from the increase in the affinity of TnC for Ca\textsuperscript{2+}, but from enhanced thin filament cooperative activation.

Dobesh et al. (2002) reported that the nH of the force–pCa curve is not correlated with \( \Delta pC_{30} \), leading them to conclude that thin filament cooperative activation plays no significant role in determining the magnitude of length-dependent activation. Consistent with this finding, we observed no significant correlation between the nH of the force–pCa curve and \( \Delta pC_{50} \) (Fig. 5). However, it is unclear to what extent the steady-state nH reflects thin filament cooperative activation. For instance, NEM-S1 or MgADP has been used to enhance thin filament cooperative activation; however, both NEM-S1 (Swartz and Moss, 1992, 2001; Fitzsimons and Moss, 1998; Fitzsimons et al., 2001a,b) and MgADP (at high concentrations: Fukuda et al., 1998, 2000) reportedly decrease the nH of the force–pCa curve, resulting presumably from enhanced recruitment of neighboring cross-bridges, especially at low Ca\textsuperscript{2+} concentrations (see Fukuda et al., 1998 for ADP contraction occurring in the absence of Ca\textsuperscript{2+}). Therefore, in the present study, we regarded the rate of rise of active force as an index of thin filament cooperative activation (as in, e.g., Swartz and Moss, 1992, 2001; Fitzsimons et al., 2001a,b), rather than the nH of the force–pCa curve.

Earlier, we discussed that at high activation states (i.e., high Ca\textsuperscript{2+} concentrations, MgADP application, or sTn reconstitution), cross-bridge recruitment upon SL elongation becomes less pronounced due to a decrease in the fraction of recruitable cross-bridges (that can potentially generate active force), resulting in the attenuation of length-dependent activation (see Fukuda et al., 2009 and references therein). The present model calculation provides a mechanistic insight into this interpretation. Namely, the SL elongation (i.e., lattice reduction)–induced increase in the probability of cross-bridge formation becomes less pronounced upon the increase in thin filament cooperative activation. This is because the acceleration of Ca\textsuperscript{2+}-dependent widening of the Gaussian distribution, of which the magnitude depends on nH\textsubscript{actin} (see Eq. 6 and Fig. S1), diminishes the lattice spacing–dependent change of the actomyosin interaction (determined by \( d\text{"a} \)). On the other hand, \( pC_{50}\text{,actin} \) had little effect on the lattice spacing dependence; i.e., length-dependent activation. Indeed, the attenuation of length-dependent activation upon sTn reconstitution was quantitatively simulated by our model (Fig. 7 A) by increasing nH\textsubscript{actin} with appropriate pC\textsubscript{30} values for both short and long SLs as a result of an increase in \( pC_{50}\text{,actin} \).

In addition, our model could quantitatively simulate the relationship of SL versus active force at various Ca\textsuperscript{2+} concentrations, converted from the force–pCa curves obtained under the control condition in Figs. 2 and 3 (i.e., shallower at high Ca\textsuperscript{2+} concentrations; Kentish et al., 1986; Fukuda et al., 2001b; see Fig. S5), emphasizing the adequacy of our model to analyze the molecular mechanism of length-dependent activation.

However, it is important to discuss limitations of the present study. First, we noted a mismatch between the experimental data and the simulated curves; namely, an increase in nH\textsubscript{actin} increased the steepness of the force–pCa curve at both SLs (compare Fig. 7 A), whereas the nH of the force–pCa curve was not increased upon sTn reconstitution or MgADP application at either SL (Table I). We consider that this mismatch reflects the limitation of the experiments with skinned myocardial fibers. For example, internal sarcomere shortening that presumably occurs during isometric contraction (Fukuda et al., 2001b)

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<th>SL (( \mu m ))</th>
<th>Passive force (mN/mm\textsuperscript{2})</th>
<th>Maximal force (mN/mm\textsuperscript{2})</th>
<th>( pC_{50} )</th>
<th>( \Delta pC_{50} )</th>
<th>nH</th>
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<td>Vehicle</td>
<td>1.9</td>
<td>55.58 ± 3.97</td>
<td>5.55 ± 0.01</td>
<td>3.72 ± 0.14</td>
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<td>2.3</td>
<td>56.80 ± 3.24</td>
<td>5.77 ± 0.01</td>
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<td>+Pimobendan</td>
<td>1.9</td>
<td>55.58 ± 3.97</td>
<td>5.55 ± 0.01</td>
<td>3.72 ± 0.14</td>
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<td>2.3</td>
<td>61.81 ± 5.92</td>
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<td>Vehicle</td>
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Data are for Fig. 6. Maximal force was obtained by activating muscle at pCa 4.5 before construction of the force–pCa curve at each SL (passive force was measured just before activation at pCa 4.5). Maximal force obtained before sTn reconstitution: 56.06 ± 3.42 and 54.25 ± 2.20 mN/mm\textsuperscript{2} for vehicle and pimobendan-treated group, respectively (\( P > 0.05 \) compared with the value obtained after sTn reconstitution). Pimobendan did not significantly change maximal force (see Fukuda et al., 2000) with and without sTn reconstitution, and it did not change any parameter in sTn-reconstituted muscles. \( ^{*}P < 0.05 \) compared with the corresponding values in the vehicle-treated group without sTn reconstitution.

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may decrease active force production by a greater magnitude at high Ca$^{2+}$ concentrations, resulting in an underestimation of the steepness of the force–pCa curve (as discussed in Fukuda et al., 2005). It should also be stressed that the mismatch reflects the limitation of the use of $t_{1/2}$ as an index of thin filament cooperative activation; namely, MgADP, Pi, or sTn reconstitution may alter the cross-bridge kinetics via a pathway that is not coupled with thin filament cooperative activation. Clearly, future studies with various techniques are needed to clarify this issue. Second, the decrease in the intermolecular distance, i.e., $d-a$ upon the addition of MgADP (earlier assumed to represent the lattice spacing modulation via cross-bridge formation in fast skeletal muscle; Shimamoto et al., 2007) enhanced length-dependent activation (Fig. S6) in contrast to the experimental result (Fig. 2). This apparent discrepancy may suggest that a change in thin filament cooperative activation has a greater impact on length-dependent activation, masking the effect of a cross-bridge–dependent lattice spacing change.

The 3-D graph obtained in the present model calculation (Fig. 7 B) suggests the role of the thin filaments in the regulation of length-dependent activation; namely, the magnitude of this phenomenon depends only slightly on the Ca$^{2+}$-binding ability of TnC, as confirmed by our experimental analysis with pimobendan (Fig. 6 and Fukuda et al., 2000), but rather strongly on the cooperativity of the thin filament on–off switching (Figs. S1 and S2). Here, it is worthwhile noting that $pCa_{50_{\text{actin}}}$ needed to be varied to quantitatively simulate the experimentally obtained relationship of $pCa_{50}$ versus $\Delta pCa_{50}$ (Fig. 7 B). This may be due to a coupling between thin filament cooperative activation and cross-bridge formation, and to the ensuing feedback effect that enhances Ca$^{2+}$ binding to TnC (Güth and Potter, 1987; Kurihara and Komukai, 1995). Therefore, the inverse relationship between $pCa_{50}$ and $\Delta pCa_{50}$ (Fig. 4 A) may be an apparent phenomenon resulting from enhanced Ca$^{2+}$ binding to TnC, coupled with acceleration of thin filament cooperative activation.

However, we admit that the present modeling is not suitable to account for the differing magnitudes of length-dependent activation in fast skeletal muscle versus slow skeletal muscle. Indeed, Konhilas et al. (2002a) reported that length-dependent activation is less in slow skeletal muscle, despite a lesser magnitude of thin filament cooperative activation. It is therefore likely that the difference in the magnitude of length-dependent activation between fast skeletal muscle and slow skeletal muscle results from factors that do not involve thin filament cooperative activation, such as isoform variance of

Figure 7. Simulation of active force development. (A) Force–pCa curves at SL 1.9 and 2.3 µm. (Left) Absolute data. (Right) Normalized data (normalized at pCa 4.5). Symbols are experimental data (the same as in Figs. 2 and 3; see Table I for absolute data). Error bars are not shown for simplicity. And solid lines are simulation results (black, control; red, sTn reconstitution). The differences between the values of the midpoint of the force–pCa curve at SL 1.9 and 2.3 µm are 0.24 and 0.14 pCa units in control and sTn-reconstituted PLV, respectively (see Table I). (B) 3-D graph showing the relationship between $n_{H_{\text{actin}}}$, $pCa_{50_{\text{actin}}}$, and $\Delta pCa_{50}$ obtained from the model analysis. Calculations were conducted at various values of $n_{H_{\text{actin}}}$ per unit of 0.5 and $pCa_{50_{\text{actin}}}$ per unit of 0.25, and, thereby, the $\Delta pCa_{50}$ values obtained from Eq. 6 (refer to Materials and methods) are plotted as the continuous mesh blue plane. Red points indicate the pairs of $n_{H_{\text{actin}}}$ and $pCa_{50_{\text{actin}}}$ that fulfill the linear relationship between $pCa_{50}$ and $\Delta pCa_{50}$ (as in Fig. 4 A).
thin filament– and thick filament–based proteins. It is an area of future research to clarify this issue by using various skeletal muscle tissues.

Considering that Ca\(^{2+}\) sensitivity of force varies depending on the type of heart disease (for review see Ohtsuki and Morimoto, 2008), it is likely that thin filament cooperative activation is altered in disease. The findings of the present study suggest that length-dependent activation is modulated via mutation occurring in the thin filaments and/or breakdown of ATP and the ensuing elevations in ADP and Pi in the vicinity of cross-bridges. Indeed, it has been reported that the Frank-Starling mechanism is depressed in skinned left ventricular muscles from patients with terminal heart failure (Schwinger et al., 1994; Brixius et al., 2003). It would be interesting to simulate, based on the 3-D state diagram (Fig. 7 B), how the Frank-Starling relation is altered by the occurrence of a mutation in a regulatory protein and/or the breakdown of ATP in various types of heart disease in various animal species, including humans.

In conclusion, thin filament cooperative activation plays a central role in the regulation of the Frank-Starling mechanism of the heart.

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Model analysis

The experimental data were fitted as follows. First, we determined the maximal amplitude of thin filament fluctuation ($\sigma_{\text{max}}$; see Figs. S1 and S2) by fitting the ratio of maximal isometric force at SL 1.9 and 2.3 μm (i.e., ~1.3 times greater at SL 2.3 μm than at 1.9 μm; see Table I in the main text). This is based on the assumption that the equilibrium of the thin filament state is fully shifted toward “on” at the saturating Ca$^{2+}$ concentration (pCa 4.5) and, accordingly, the lattice spacing predominates over thin filament fluctuation in producing the length-dependent activation. Indeed, the SL dependence of maximal Ca$^{2+}$-activated force is in agreement with what was observed in previous experiments by various groups, including us (Fukuda et al., 2000, 2001, 2003, 2005; Terui et al., 2008). For quantitative simulation of the force–pCa curves at SL 1.9 and 2.3 μm, we set the value of the parameter $\sigma_{\text{max}}$ at 21 nm, similar to the value used in the previous model analysis (Ishiwata and Oosawa, 1974). At larger values of $\sigma_{\text{max}}$ the ratio of the calculated maximal force at these SLs was too small, and at smaller values, the force–pCa curve exhibited higher than normal steepness compared with the experimental data (not depicted).

Second, we determined the average position of the myosin heads, $a$, under a fixed $\sigma_{\text{max}}$ value (i.e., 21 nm) so as to simulate the force–SL curves at various Ca$^{2+}$ concentrations, which were experimentally obtained. As documented in our previous studies (Ishiwata and Oosawa, 1974; Shimamoto et al., 2007), not only the lattice spacing, $d$, but also the average position of the myosin heads, $a$, determines the interaction probability of actin and myosin at various SLs (compare Eq. 6 in the main text). The value, $a$, was set at 24 nm, with which the force–SL curves were quantitatively simulated at various pCa values (Fig. S5). With SL >2.8 μm, the calculated forces were found to be higher at low Ca$^{2+}$ concentrations than at high Ca$^{2+}$ concentrations (not depicted) because the actomyosin interaction starts to occur at low Ca$^{2+}$ concentrations due to overly reduced lattice spacing (which is smaller than the width of the Gaussian distribution; compare Fig. S1). However, given the fact that adult mammalian ventricular muscle cannot be stretched beyond SL ~2.4 μm under normal physiological conditions, because of irreversible damage of the muscle structure (Weiwad et al., 2000), we disregarded this model limitation.

Third, and finally, we varied $n_{\text{H,actin}}$ between 1 and 5 and $p\text{Ca}_{\text{50,actin}}$ between 5.0 and 6.0 in the XY plane, and plotted the $\Delta p\text{Ca}_{\text{50}}$ value obtained at each parameter pair on the Z axis (that was calculated from the simulated force–pCa curves at SL 1.9 and 2.3 μm). Thereby, the 3-D graph was constructed (Fig. 7 B).

To summarize our model, active isometric force was determined by multiplying the overlap length of the thick and thin filament with the probability of actomyosin interaction per unit length, and the probability was assumed to change depending on the Ca$^{2+}$ concentration and lattice spacing. Again, because of the small change in the overlap between SL 1.9 and 2.3 μm, the overlap length, i.e., $1/2(L_0-L)$, was assumed to be constant. For the quantitative simulation of the experimental data, the value of $F_0$ was chosen to be 45 and 36, for control and s1Tn-reconstituted PLV, respectively (see Eq. 6).
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Figure S2. Graphs showing the characteristics of our model. (A) Plots of pCa versus σₘ. σₘ was assumed to be increased as a function of pCa, according to the Hill equation (Eq. 4 in the main text). In this case, nₜₐₜₐ₃ and pCaₐₜₐ₃ were set at 1.51 and 5.44 for control PLV, and at 2.54 and 5.72 for sTn-reconstituted PLV, respectively. (B) Effect of varying nₜₐₜₐ₃ on the midpoint of the force–pCa curve, i.e., pCaₐₜₐ₃, at a fixed value of pCaₐₜₐ₃, i.e., 5.44, a value used for the simulation of the force–pCa curves for control PLV. pCaₐₜₐ₃ was decreased at both SLs, accompanied by diminution of ΔpCaₐₜₐ₃. The filled circle and triangle indicate the midpoint of the force–pCa curve for SL 2.3 μm, respectively, at the nₜₐₜₐ₃ value of 1.51. (C) Same as in B, but nₜₐₜₐ₃ was varied at pCaₐₜₐ₃, 5.72, a value used for the simulation of the force–pCa curves for sTn-reconstituted PLV. The filled circle and triangle indicate the midpoint of the force–pCa curve (see Fig. 7 A) for SL 1.9 and 2.3 μm, respectively, at the nₜₐₜₐ₃ value of 2.54.
Figure S3. Cross-sectional view of the meshed blue plane in the 3-D graph (Fig. 7B) at a fixed value of $n_{H_{\text{actin}}}$ (A) and $pCa_{50_{\text{actin}}}$ (B). It is clearly seen that $\Delta pCa_{50}$ is barely changed with $pCa_{50_{\text{actin}}}$ but markedly decreased upon the increase in $n_{H_{\text{actin}}}$ (especially up to 3.0).

Figure S4. Effect of MgADP on $k_{tr}$ in PLV. $k_{tr}$ was measured based on our previous study (Terui et al., 2008). (A) Typical recording showing active force redevelopment in the absence (left) and presence (right) of 3 mM MgADP at maximal ($pCa_{4.5}$) and half-maximal (i.e., $pCa_{50}$; see Table I) activations. The same preparation was used throughout the experiment for each graph ($k_{tr}$ measured first at half-maximal activation and then at $pCa_{4.5}$). SL, 1.9 μm. Maximal force value at steady state: 31.46 (39.02) and 63.90 (65.94) mN/mm² at half-maximal and maximal activation, respectively. Maximal $k_{tr}$ values: 0.86 ± 0.19 and 0.76 ± 0.14 s$^{-1}$ in the absence and presence of MgADP, respectively. Maximal force values were 63.23 ± 11.65 and 53.43 ± 7.02 mN/mm² in the absence and presence of MgADP, respectively. $n = 4$. (B) Force-$k_{tr}$ relationship (see Terui et al., 2008). At similar values of relative force ($P > 0.05$), $k_{tr}$ was significantly lower in the presence of MgADP.
Figure S5. Relationship between SL and active force at various Ca^{2+} concentrations in PLV (control condition) simulated by our model. The SL-active force relationship became shallow upon the increase in the Ca^{2+} concentration, consistent with previously obtained experimental results with skinned as well as intact ventricular preparations (Kentish et al., 1986; Fukuda et al., 2001). Circles with solid lines, simulation results; triangles with dashed lines, experimental results (compare Figs. 2A and 3A, and Table I). Note that the present experimental results were quantitatively simulated by our model.

Figure S6. Model simulation showing the effect of an increase in the average length of myosin heads, \(a\), on Ca^{2+} sensitivity and the length-dependent activation in control PLV. The parameter \(a\) was varied between 22 and 25 nm (\(a=24\) nm in the current modeling; see the main text). \(p_{\text{Ca}_{50}}\) and \(p_{\text{Ca}_{50}/\text{actin}}\) were set at 1.51 and 5.44, respectively (compare Fig. S2B). Ca^{2+} sensitivity of force was found to be increased upon the increase in \(a\) at both SL 1.9 and 2.5 \(\mu\)m (A), accompanied by enhancement of length-dependent activation (B).