Basolateral K Channels
in an Insect Epithelium

Channel Density, Conductance,
and Block by Barium

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ABSTRACT K channels in the basolateral membrane of insect hindgut were studied using current fluctuation analysis and microelectrodes. Locust recta were mounted in Ussing-type chambers containing Cl-free saline and cyclic AMP (cAMP). A transepithelial K current was induced by raising serosal [K] under short-circuit conditions. Adding Ba to the mucosal (luminal) side under these conditions had no effect; however, serosal Ba reversibly inhibited the short-circuit current (Isc), increased transepithelial resistance (Rt), and added a Lorentzian component to power density spectra of the I.. A nonlinear relationship between corner frequency and serosal [Ba] was observed, which suggests that the rate constant for Ba association with basolateral channels increased as [Ba] was elevated. Microelectrode experiments revealed that the basolateral membrane hyperpolarized when Ba was added: this change in membrane potential could explain the nonlinearity of the 2πf vs. [Ba] relationship if external Ba sensed about three-quarters of the basolateral membrane field. Conventional microelectrodes were used to determine the correspondence between transepithelially measured current noise and basolateral membrane conductance fluctuations, and ion-sensitive microelectrodes were used to measure intracellular K activity (aK). From the relationship between the net electrochemical potential for K across the basolateral membrane and the single channel current calculated from noise analysis, we estimate that the conductance of basolateral K channels is ~60 pS, and that there are ~180 million channels per square centimeter of tissue area.
INTRODUCTION

Many epithelial cells in vertebrate and invertebrate animals have a large K conductance in their basolateral membrane that allows K ions entering through the Na/K pump to return to the serosal side (Koefoed-Johnsen and Ussing, 1958; reviewed by Lewis et al., 1984b). Basolateral K conductance influences the electrical potential across the apical membrane because of the finite conductance of the paracellular shunt, and is therefore important in determining flux through apical ion channels and electrogenic cotransporters (e.g., Morgunov and Bouapaep, 1984; Lapointe et al., 1984). In addition, basolateral K conductance has been implicated in the regulation of cell volume in several epithelia (Germann et al., 1984; Lau et al., 1984; Lewis et al., 1984a). While inhibition of basolateral K conductance by millimolar levels of Ba is well established in many epithelia, the kinetics and voltage dependence of Ba block have not been examined. In fact, relatively little is known about basolateral K conductance at the single channel level in absorptive epithelia, although exocrine glands have been studied using the patch-clamp technique (e.g., Maruyama et al., 1983a, b).

Current fluctuation analysis has been used extensively for studies of epithelial ion channels. By assuming a simple model for blocker interactions, noise analysis has allowed the calculation of single channel currents and the densities of apical Na and K channels (e.g., Lindemann and Van Driessche, 1977; Van Driessche and Zeiske, 1980). However, this approach has been less useful for studying basolateral channels because many preparations have a high apical membrane resistance, which leads to attenuation of transepithelial current noise relative to conductance fluctuations at the basolateral membrane. This problem can be minimized by using ionophores such as nystatin to lower apical membrane resistance artificially. For example, spontaneous current noise from the basolateral membrane has been measured under these conditions in rabbit colon (Wills et al., 1982; see Wills, 1984) and amphibian skin (Van Driessche et al., 1982). Nonetheless, it has not been possible so far to estimate basolateral single channel currents or their conductance, nor has the number of channels per unit tissue area been assessed in these epithelia.

One preparation that is uniquely suited for studying basolateral K channels by this technique is the hindgut of the desert locust. The locust renal system consists of Malpighian tubules, which actively secrete a K-rich primary urine (140–165 mM K) into the hindgut, and rectal pads that selectively reabsorb water, salts, and nutrients (see Phillips, 1981). Flat-sheet cable analysis has shown that the locust rectum is a tight epithelium, and that most of the K secreted by Malpighian tubules is recovered passively in the rectum through conductive pathways in the apical and basolateral membranes (Hanrahan and Phillips, 1984b, c). Furthermore, apical K conductance is under hormonal control in locust rectum, and can be stimulated fivefold by cAMP in vitro without affecting the K conductance of the basolateral membrane (Hanrahan and Phillips, 1984b).

In this paper, we exploit these unusual properties of insect hindgut to study basolateral K channels. cAMP-stimulated recta are short-circuited in the presence of a transepithelial K gradient, and Ba-induced fluctuations in $I_K$ are analyzed. We estimate the kinetics of block and the single channel current from transep-
the epithelial noise, and use microelectrode data collected under identical conditions to correct for the attenuating effects of the apical membrane. The corrected single channel current is then used to calculate channel density and to determine channel conductance from Ba-induced variations in K driving force.

Some of these data were reported previously in abstract form (Hanrahan et al., 1983).

METHODS

Preparations

To measure transepithelial current noise, recta from adult male desert locusts (Schistocerca gregaria, Forskål) were dissected and mounted in Ussing-type chambers (Lewis and Diamond, 1976) modified to expose 0.239 cm² of tissue. Each half-chamber had a volume of 15 ml and was vigorously stirred and oxygenated except when records were actually being collected for analysis. Tissues were allowed to equilibrate for 30 min in normal saline containing (mM): 70 Na, 10 K, 10 Mg, 5 Ca, 110 Cl, 10 glucose, 100 sucrose, 10 proline, and 5 glutamine at pH 7.2. Based on previous studies of this tissue (Hanrahan and Phillips, 1984a–c), we optimized the experimental conditions for studying transcellular K current (see Fig. 1, A and B). First, Cl was replaced with gluconate on both sides to abolish electrogenic Cl absorption and basolateral Cl conductance. Second, 1 mM cAMP was added to the serosal side to stimulate apical membrane K conductance. Finally, a transcellular K current was induced by replacing serosal Na with K.

Under these short-circuit conditions, we would expect the chemical gradient across the basolateral membrane to be virtually eliminated and intracellular potential to result from the outward K gradient at the apical membrane. A serosal-to-mucosal (s → m) rather than a mucosal-to-serosal (m → s) K gradient was chosen because there is evidence that apical K permeability declines in the presence of high mucosal [K] (Hanrahan and Phillips, 1984a). To induce current fluctuations, we added 0–20 mM Ba(NO₃)₂ to the serosal side from a concentrated aqueous stock solution (200 mM). The same volume of 400 mM NaNO₃ was added to the mucosal side.

For intracellular recording, we used a microelectrode chamber that has been described in detail elsewhere (Hanrahan et al., 1984b; Hanrahan and Phillips, 1984a). Aside from the different chambers used, other experimental conditions during intracellular measurements were identical to those described above for noise analysis. Both chamber designs have been shown to cause minimal edge damage. All experiments were done at room temperature (~22°C).

Electrical Methods

A low-noise voltage clamp similar to the one described by Van Driessche and Lindemann (1978) was used to measure short-circuit current. Data collection and analyses were similar to those described previously (Loo et al., 1983; Lewis et al., 1984c; Wills et al., 1984). Briefly, electrodes (voltage-measuring and current-passing) were Ag/AgCl wires imbedded in 4% agar/1 M KCl. Voltage-sensing agar bridges were positioned as close to the tissue as possible and ~30° normal to the plane of the epithelium using micromanipulators. $I_c$ was recorded after high-pass-filtering at 0.03 Hz (113, EG&G Princeton Applied Research, Princeton, NJ), amplified, low-pass-filtered through a 20-pole Cauer-Elliptic filter set at 100 Hz to prevent aliasing (LP120, Unigon Industrial Inc., Mount Vernon, NY), and sampled every 5 ms. The resulting 32,768 points were fast-Fourier-transformed in blocks of 2,048, beginning after every 512th point. The final power spectra, which represent an average of 61 transforms obtained from 16 overlapping data blocks, were fitted using a
nonlinear curve-fitting routine based on the derivative-free Levenberg-Marquardt algorithm of Brown and Dennis (1972). Following Van Driessche and Zeiske (1980), spectra were fitted as the sum of linear and Lorentzian-type functions:

$$S_i(f) = A/f^\beta + S_i(0)/(1 + (ff_0)^\gamma),$$

where $A$ is defined as the power at 1 Hz, $\beta$ is the slope of the linear component, $S_i(0)$ is

**Figure 1.** (A) Diagram of KCl transport across locust rectum bathed in normal saline containing 110 mM Na, 10 mM K, and 110 mM Cl. The peptide "chloride transport-stimulating hormone" (CTSH) elevates intracellular cAMP (Spring and Phillips, 1980). Cyclic AMP stimulates (a) apical K conductance, (b) active Cl influx, and (c) basolateral membrane Cl conductance (Hanrahan and Phillips, 1984). (B) Experimental conditions used to study basolateral K conductance. CI was replaced with gluconate on both sides. Cyclic AMP (1 mM) was used to stimulate apical K conductance. $I_{Kb}$ was induced by raising serosal [K] from 10 to 80 mM. These conditions were designed to optimize basolateral membrane voltage clamp when $V_t$ was short-circuited. (C) Equivalent circuit of locust rectum under the experimental conditions shown in B, including the voltage clamp. $E_a$ and $E_b$ are the emfs of apical and basolateral membranes, respectively, $R_e$ is the resistance of the apical membrane, and $R_b$ is the resistance of the basolateral membrane, which depends on $[Ba]$. $g_{Kb}$ is the single channel conductance, $M$ is the number of channels, and $P_o$ is the probability that a channel is unblocked. For simplicity, the high-resistance paracellular pathway and its emf have not been included: under short-circuit ($I_m$) conditions, paracellular current would not flow through $R_b$ and can be subtracted when calculating $I_{Bb}$ (see text). $V_b$ is the basolateral membrane potential, which under short-circuit conditions is identical to the apical membrane potential.
the low-frequency plateau of the Lorentzian component, and \( f_c \) is the corner frequency. The macroscopic \( I_K \) was monitored on a strip-chart recorder and digitized by a 12-bit A/D converter interfaced to a microcomputer (NorthStar Horizon II).

Single-barreled microelectrodes were pulled (PP-88, Narishige Scientific Instruments Laboratory, Tokyo, Japan) from borosilicate tubing containing a solid fiber (1.0 mm o.d., Frederick Haer & Co., Brunswick, ME). Conventional microelectrodes were filled with 0.5 M KCl and had resistances of 16.6 ± 1.5 MΩ. The tips of K-selective microelectrodes were silanized, filled with resin (477317, Corning Glassworks, Corning, NY), and back-filled with 0.5 M KCl as previously described (Hanrahan et al., 1984b; Hanrahan and Phillips, 1984b). Ion-sensitive microelectrodes were calibrated in solutions containing 1, 10, 50, and 100 mM KCl immediately before or after making a series of ~12 impalements. The electrodes used in the present experiments had a mean slope of 53.24 ± 0.58 mV/decade change in K activity (± SE, 33 electrodes), as calculated from the modified Debye-Hückel theory (Robinson and Stokes, 1959). No correction was made for interference by intracellular Na (\( a_{\text{Na}} \)) because electrodes were relatively selective for K over Na according to the separate solution method in which electrode emfs are compared in pure solutions (selectivity for K over Na = 62.4 ± 2.2), and because \( a_{\text{Na}} \) in locust rectal cells bathed with normal saline (~8 mM; Hanrahan and Phillips, 1984b) is normally too low to affect K activity measurements significantly. Moreover, \( a_{\text{Na}} \) is probably reduced further under the present high-K conditions.

Transepithelial potential (\( V_e \)) was arbitrarily clamped to 7 mV for 110 ms and the resulting deflections in apical (\( V_a \)) and basolateral membrane potentials (\( V_b \)) were recorded. We estimated \( \alpha \), the ratio of apical-to-basolateral membrane resistance (\( R_a/R_b \)), according to

\[
\alpha = \frac{\Delta V_a}{\Delta V_b} \approx \frac{R_a}{R_b}.
\]

The voltage drop caused by saline resistance between the voltage-sensing agar bridges and epithelium was subtracted when calculating \( \alpha \), \( R_a \), and the net driving force for K across the basolateral membrane. The effect of mucosal saline resistance on \( \Delta V_a \) was very small (<0.5 mV, barely detectable in Fig. 7 as deflections in the trace between impalements), although the saline resistance was somewhat larger on the serosal side and resulted in deflections of ±1.6 mV.

Net electrochemical potential for K across the basolateral membrane (\( \Delta \mu_K/F \)) was calculated as

\[
\Delta \mu_K/F = (RT/F) \ln \left( \frac{a_{\text{K}}}{a_{\text{K}^*}} \right) + V_b,
\]

where \( a_{\text{K}} \) and \( a_{\text{K}^*} \) are intracellular and serosal K activities, respectively, and \( R, T, F \), and \( V_b \) have their usual meanings. Statistical significance was determined using unpaired t tests.

**Equivalent Circuit Model**

Fig. 1C shows an equivalent circuit model for locust rectum, which is useful when interpreting the data presented below. We note the following features. (a) Apical membrane resistance (\( R_a \)) is minimized by adding cAMP (Hanrahan and Phillips, 1984b). (b) The [K] gradient from cell-to-mucosa (measured in Results) is largely responsible for \( E_a \) because apical Na conductance is negligible, and because both compartments are nominally Cl-free. (c) Junctional resistance (\( R_j \)) and the diffusion potential within the paracellular pathway (\( E_j \)) caused by s → m K gradient and m → s Na gradient) are omitted from this simplified circuit. Locust rectum is a very tight epithelium during cAMP stimulation.
according to cable analysis, so that little current would flow paracellularly (Hanrahan and Phillips, 1984b). Furthermore, paracellular current is excluded from subsequent calculations by subtracting the Ba-insensitive current (see below). (d) A basolateral membrane emf (approximately -5 mV) is included because the basolateral K gradient was not completely abolished under our experimental conditions (see Results). (e) Basolateral membrane conductance is shown as a variable resistor that depends on the Ba concentration. (f) The saline resistance between voltage-sensing agar bridges and the epithelium is not included; however, series resistance was minimized during noise experiments by positioning the agar bridges very close to the tissue using micromanipulators, and corrected for in microelectrode experiments as described above.

RESULTS

Effects of Serosal Ba on $I_w$

When recta were mounted in Ussing-type chambers under control, open-circuit conditions, $V_t$ decreased during the first hour from 16.6 ± 2.5 to 4.7 ± 1.8 mV ($\bar{x} \pm$ SE, 11 recta), results comparable to those obtained previously using a similar saline (Hanrahan and Phillips, 1984a). Fig. 2A shows the effect on $V_t$ of increasing...

![Figure 2](image-url)

**Figure 2.** (A) Effect of raising serosal K concentration from 10 to 80 mM under open-circuit conditions ($V_t$, referenced to the serosal solution). Both sides were nominally Cl-free (gluconate-substituted) and 1 mM cAMP was present on the serosal side. Brief deflections in the trace result from constant-current pulses (1.8 $\mu$A), which were used to monitor tissue resistance. Preparation LOC0122. (B) Effect of serosal Ba on short-circuit current ($I_s$). Most $I_s$ was Ba sensitive. Deflections in $I_s$ became smaller when Ba was added, which indicated an increase in transepithelial resistance ($R_t$). Preparation LOC0907. (C) Effect of Ba on the macroscopic K current ($I_{Kw}$; Ba-insensitive current subtracted) and $R_t$. Mean ± SE.
serosal K concentration from 10 to 80 mM in the presence of 1 mM cAMP under Cl-free conditions. When $V_i$ was clamped to 0 mV, the mean $I_{sc}$ was variable (53.9 ± 11.0 μA cm$^{-2}$); however, most of the current was abolished by adding Ba to the serosal side (Fig. 2B). We did not observe a gradual increase in $I_{sc}$ after prolonged exposure to Ba like that described in frog skin (Van Driessche

![Graph showing the effect of Ba concentration on $I_{sc}$ over time and [Ba] concentration on $I_{sc}$ and $R_i$.)

FIGURE 2, B and C
and Zeiske, 1980). However, there was often a small decline in transepithelial resistance when serosal \([\text{Ba(NO}_3\text{)}_2]\) was elevated beyond 10 mM, probably a consequence of nitrate conductance. To minimize this problem, only data obtained at low Ba concentrations (\(\leq 10\) mM) were used when calculating the blocking kinetics and channel properties.

We obtained \(I_{\text{KBa}}\), the macroscopic current flowing through Ba-sensitive channels, at each Ba concentration by subtracting the \(I_{\text{c}}\) observed with 20 mM or 13.1 mM Ba present on the serosal side from that measured; i.e.,

\[
I_{\text{KBa}} = I_{\text{c}} - I_R,
\]

where \(I_R\) is the residual current remaining when all the Ba-sensitive channels are blocked. Fig. 2C shows the decline in \(I_{\text{KBa}}\) and the increase in \(R_c\) caused by adding Ba to the serosal side. Half-maximal inhibition of \(I_{\text{KBa}}\) (i.e., the macroscopic \(K_i\)) appears to be in the range of 2.0–3.0 mM Ba, which is similar to that reported for external Ba block in other cells (e.g., squid axon; Armstrong et al., 1982).

**Ensemble Analysis of Ba-induced Current Fluctuations**

Adding millimolar levels of barium to the serosal (high-K) solution induced fluctuations in the short-circuit current (5 mM Ba; Fig. 3A), which decreased at very high Ba concentrations (20 mM). In the complete absence of Ba ("control"; Fig. 3B) or when Ba was added only to the mucosal side, the power spectral density of the \(I_{\text{sc}}\) noise could not be fitted with a Lorentzian component. By contrast, serosal addition of Ba induced a Lorentzian component in the spectrum (5 mM Ba; Fig. 3B). The low-frequency plateau value \(S_0\) decreased, and the corner frequency \((f_c)\) increased as the serosal Ba concentration was raised stepwise (Fig. 3C).

Inhibition of the macroscopic current by Ba was completely reversible; therefore, we assumed the simplest scheme for channel (R) block, which has also been used by others (e.g., Standen and Stanfield, 1978; Armstrong and Taylor, 1980; Eaton and Brodwick, 1980; Van Driessche and Goegelein, 1980; Vergara and Latorre, 1983; DeCoursey and Hutter, 1984):

\[
R + \text{Ba} \rightarrow \frac{k_{01}}{k_{10}} (R\text{Ba}).
\]

As demonstrated previously for amiloride-induced noise (Lindemann and Van Driessche, 1977), if the number of channels at the membrane surface is negligible compared with the concentration of Ba, the overall reaction rate \((1/\tau)\) at equilibrium is

\[
1/\tau = [\text{Ba}]k_{01} + k_{10} = 2\pi f_c,
\]

and the blocking reaction is pseudo-first order, so that plotting \(2\pi f_c\) vs. \([\text{blocker}]\) yields a straight line with slope = \(k_{01}\) (with units \(\text{mM}^{-1} \text{s}^{-1}\)) and intercept = \(k_{10}\) (in \(\text{s}^{-1}\)). Fig. 4 shows the relationship between \(2\pi f_c\) and serosal \([\text{Ba}]\) obtained for locust rectum. The correlation coefficient determined by linear regression was quite high \((r^2 = 0.9965)\), although we noticed that the value of \(2\pi f_c\) at 3.5 mM Ba was always below the regression line.
FIGURE 3. Ba-induced fluctuations in $I_{sc}$. Ba was added to the serosal side in the presence of a serosal-to-mucosal K gradient and 1 mM cAMP. (A) Fluctuations in short-circuit current were measured with 0, 5, or 20 mM Ba on the serosal side. (B) Power spectra obtained in the absence of Ba (O) could not be fitted with a Lorentzian-type function; however, a Lorentzian component was observed after Ba was added to the serosal side (+). (C) Effect of increasing serosal [Ba] on power spectra. Best fits to Eq. 1 are shown at each concentration for one preparation (see text). Preparation LOC0122.
Effect of Ba on the Apparent Single Channel Current

Applying the simplest "on-off" blocking scheme and a two-state kinetic model (e.g., Neher and Stevens, 1977), the low-frequency limit of \( S(f) \) is

\[
S(0) = \frac{i_{KBa} M P_o [Ba]}{f_c^2},
\]

where \( i_{KBa} \) is the current flowing through a single (unblocked) channel in the basolateral membrane (see Fig. 1C), \( M \) is the number of channels, \( P_o \) is the probability that a single channel is unblocked at a particular concentration of Ba, and \( f_c \) is the corner frequency of the Lorentzian component at that concentration.

Re-examination of the \( 2\pi f_c [Ba] \) Plot

The apparent single channel currents (\( i_{KBa} \); reasons for the use of "apparent" will become clear in a later section) were calculated for each [Ba] by rearranging Eq. 6 to give

\[
i_{KBa} = \frac{i(0) P_o f_c^2}{I_{KBa} [Ba] k_{01}},
\]

and inserting \( I_K, S(0), f_c, \) and the \( k_{01} \) obtained from individual preparations over the entire [Ba] range (0.5–10 mM). Although the single channel current should not depend on Ba concentration, Fig. 5 shows that \( i_{KBa} \) was approximately eightfold larger at 10 mM [Ba] compared with that at 0.5 mM Ba. The increase in \( i_{KBa} \) with [Ba] might be explained by several factors: (a) the \( 2\pi f_c [Ba] \) plot may not be linear, as evidenced by the low value at 3.5 mM in Fig. 4, (b) attenuation of basolateral membrane current noise by the apical membrane may be reduced at high Ba concentrations, and (c) the driving force for K entry may be enhanced by serosal [Ba].

Re-examination of the \( 2\pi f_c [Ba] \) Plot

When data from those five preparations in which \( f_c \) was determined at each of the lower Ba concentrations (i.e., 0.5, 1.0, 3.5, and 5.0 mM) were replotted on
FIGURE 5. Apparent single channel currents ($i_{K_{Ba}}$) calculated for individual preparations at each [Ba] using Eq. 7. $i_{K_{Ba}}$ was about eightfold larger at 10 mM Ba than at 0.5 mM Ba. Mean ± SE; see also Table I.

an expanded scale, it became obvious that the relationship between $2\pi f_c$ and [Ba] is concave upwards (Fig. 6). If one assumes that Ba interactions in locust K channels are pseudo-first order (as is the case in other preparations: e.g., Van Driessche and Gogelein, 1980; Vergara and Latorre, 1983), then a tangent to the curve at low concentrations would indicate $k_{o1} = 0.868$ mM$^{-1}$ s$^{-1}$ and $k_{o10} = 0.712$ s$^{-1}$. Clearly, this is only an approximation and should overestimate $k_{o1}$ and underestimate $k_{o10}$; nevertheless, the shape of the curve suggests that these errors would be small, and changes in membrane potential are also minimal over this range (see below).

The slope of the $2\pi f_c$-[Ba] relation was much steeper at high Ba concentrations (5–10 mM; Fig. 4), which is consistent with a value for $k_{o1}$ of $\sim$1.94 mM$^{-1}$ s$^{-1}$.

FIGURE 6. Data from five recta in which $f_c$ was obtained in the presence of 0.5, 1.0, 3.5, and 5.0 mM Ba (○: standard error at 0.5 mM Ba was smaller than the symbol). The solid line indicates $k_{o1} = 0.868$ mM$^{-1}$ s$^{-1}$ and $k_{o10} = 0.712$ s$^{-1}$, calculated for a two-state model in which the on and off rate constants are not affected by [Ba]. The open circles were calculated assuming that $k_{o1}$ depends on voltage (see Table I), with $k_{o1}(0) = 0.378$ mM$^{-1}$ s$^{-1}$ and $\gamma = 0.76$. Mean ± SE.
Although, in principle, it should be possible to calculate $k_{10}$ at high [Ba], the small value of $k_{10}$ and the relatively large extrapolation required to reach the ordinate made this impossible and in fact yielded a small negative value ($-0.16 \text{ s}^{-1}$). We believe this is due to the fact that $2\pi f_z$ is dominated by the association rate constant at high blocker concentrations (for example, $k_{10}$ would contribute only ~5% of the $2\pi f_z$ at 10 mM Ba, according to Eq. 5b and Fig. 4). As a consequence, we focus on the Ba association rate constant in the remainder of this paper, reasoning that the large contribution of the on rate constant would make $2\pi f_z$ sensitive to changes in $k_{01}$.

**Effect of Ba on Intracellular Potential**

It is well known that Ba is a voltage-dependent blocker in other preparations. The type of curvature in Fig. 6 would be expected if the basolateral membrane potential ($V_b$) hyperpolarized as the Ba concentration was elevated because this would increase the association rate constant (i.e., the slope). To determine whether Ba addition caused significant hyperpolarization, intracellular potential was measured under the same conditions present during noise measurements (i.e., short-circuited, 1 mM cAMP, transepithelial K gradient, and absence of Cl). Microelectrodes were referenced to either the mucosal or serosal agar bridge and this could be switched while recording from one cell (top of Fig. 7) or between successive impalements (bottom of Fig. 7). Serosal Ba hyperpolarized the basolateral membrane in every preparation, with $V_b$ increasing by 20.7 ± 4.5 mV ($P < 0.01$) from -13.8 ± 1.5 to -36.1 ± 12.8 when 10 mM Ba was added (Table I and Fig. 8).

We calculated the voltage dependence of block on the assumption that the curvature in the $2\pi f_z$-[Ba] relationship results from voltage-dependent association, and that the simplest (single-site, two-barrier) model adequately describes Ba interactions with the channel (Woodhull, 1973). The on rate constant for Ba block ($k_{01}$) as a function of voltage is then:

$$k_{01}(V) = k_{01}(0)\exp(-zF\gamma V/RT),$$

where $k_{01}(0)$ is the association rate constant at $V_b = 0$ mV (in mM$^{-1}$ s$^{-1}$), $\gamma$ is the fraction of the membrane field sensed by the barrier to Ba entry, and $z$, $F$, $R$, and $T$ have their usual meanings. $k_{01}(0)$ and $\gamma$ were calculated using Eq. 8 by inserting $V_b$ measured with 1 mM Ba ($-13.8$ mV; “low potential”) or 10 mM Ba ($-27.1$ mV; “high potential”), and association rate constants for low and high membrane potentials ($k_{01} = 0.868$ mM$^{-1}$ s$^{-1}$ at 0.5-1.0 mM Ba; $k_{01} = 1.94$ mM$^{-1}$ s$^{-1}$ at 5.0-10.0 mM Ba). The electrical distance of the barrier ($\gamma$) was estimated to be 0.76 of the way through the membrane field, and $k_{01}(0)$ was 0.378 mM$^{-1}$ s$^{-1}$. The values of $2\pi f_z$ calculated using these estimates for $k_{01}$ and $\gamma$ are similar to those observed, although we assumed for simplicity that $k_{10}$ is voltage independent when we made the calculation (broken line, Fig. 6). The agreement might still be improved slightly if $k_{10}$ decreased with hyperpolarization.

**Effect of Ba on the Attenuation of Basolateral Current Noise**

Short-circuiting the epithelium abolishes $V_n$, but it does not perfectly voltage-clamp the basolateral membrane, which contains the fluctuating channels. Con-
sequently, transepithelial current noise appears attenuated compared with basolateral conductance fluctuations by some amount that depends on the relative impedance of the apical and basolateral membranes (Van Driessche and Gogelein, 1980; Lindemann and DeFelice, 1981; Wills et al., 1984). When series resistance is much less than \( R_j \), transepithelially measured spectral density, \( S_t(f)_{\text{meas}} \), is related to the spectral density under ideal conditions, \( S_t(f) \), by

\[
S_t(f)_{\text{meas}} = S_t(f) |H(f)|^2,
\]

where

\[
|H(f)| = \sqrt{(Re[H(f)])^2 + (Im[H(f)])^2}
\]

and \( |H(f)| \) is the modulus of the frequency response function \( H(f) \). In principle, the entire noise spectrum should be corrected for this attenuation because it is frequency dependent and might affect estimates of \( f_c \) (Lindemann and DeFelice, 1981). However, Ba-induced noise occurs at very low frequencies in this preparation, where membrane capacitance would have little effect. Furthermore, spectra did not show any peaks that would suggest impedance artifacts (Van Driessche and Gogelein, 1980; Van Driessche and Gullentops, 1982). For these reasons, it was sufficient to determine \( H(f) \) at the low-frequency limit [i.e.,
the plateau value, $S_0(0)$ as

$$H(0) = R_b/(R_a + R_b) \approx 1/(1 + \alpha),$$

where $\alpha$ was measured at 0, 1, 5, and 10 mM [Ba] as described in the Methods.

Table I shows that the ratio of apical-to-basolateral membrane resistance ($\alpha$) declined by ~90% when 10 mM Ba was added to the serosal side, which is consistent with inhibition of most of the basolateral K conductance (Table I). The fractional resistance of the basolateral membrane $[1/(1 + \alpha)]$ increased from 0.40 to 0.87 as Ba was added, and the block was half-maximal at ~3 mM Ba, which is similar to that observed for inhibition of the macroscopic current during noise experiments. It is clear that adding Ba improves the voltage clamp of the

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<tr>
<td>Effects of Serosal Ba Addition on Intracellular Potential ($V_b$), Intracellular K Activity ($a_k$), Ratio of Apical to Basolateral Membrane Resistances ($\alpha$), and Net Electrochemical Potential for K at the Basolateral Membrane ($\Delta \mu_k/F$)</td>
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<tr>
<td>$\Delta \mu_k/F$ (mV)</td>
<td>-8.4±8.2</td>
<td>-13.6±4.8</td>
<td>-18.4±5.7</td>
<td>-25.8±4.7</td>
</tr>
</tbody>
</table>

$x \pm SE$ (n = number of tissues).
basolateral membrane, thereby reducing attenuation of the basolateral current fluctuations. Therefore, it is necessary to correct for this effect when calculating \( i_{KBa} \) by using the following equation:

\[
i_{KBa} = S_0(0)_{\text{meas}} \alpha^2 f_c^2 (1 + \alpha)^2 / I_{KBa}[Ba]k_{01},
\]

where \( S_0(0)_{\text{meas}} \) is the plateau value for the spectrum of \( I_x \) fluctuations.

### Effect of \( Ba \) on Intracellular K Activity (\( a_K \)) and Net Electrochemical Potential (\( \Delta \mu_K^i/F \))

It is already clear that serosal \( Ba \) addition under these conditions causes \( V_b \) to hyperpolarize; however, it was also necessary to measure intracellular K activity at each \( [Ba] \), so that the net driving force for K could be determined. Table I shows that \( Ba \) had little effect on intracellular K activity (\( a_K \)). It seems likely that intracellular K was maintained by basolateral Na/K-ATPase, so that only minor changes in the net driving force for K entry through the basolateral membrane were attributable to the K gradient (for biochemical and physiological evidence for this enzyme in locust rectum, see Phillips, 1981).

### Single Channel i-V Relation

The fact that \( Ba \) addition alters \( i_{KBa} \) and \( \Delta \mu_K^i/F \) enabled us to estimate single channel conductance from the relationship between K current and net driving force. Values for \( k_{01} \) at low (0.5–1.0 mM) and high (5.0–10.0 mM) \( Ba \) concentrations and \( \alpha \), measured at 1 and 10 mM \( Ba \), were inserted into Eq. 12 to obtain \( i_{KBa} \) (the single channel current corrected for changes in \( k_{01} \) and attenuation). Then \( i_{KBa} \) was plotted against the (measured) net driving force for K at 1.0 mM and 10.0 mM [Ba] as shown in Fig. 9. Chord conductance (\( g_{KBa} \)) was estimated from the relationship

\[
g_{KBa} = i_{KBa} / (\Delta \mu_K^i/F).
\]

Over the narrow range of \( \Delta \mu_K^i/F \) observed in this study, \( g_{KBa} \) was calculated to be ~60 pS. Note that this is only a first estimate because of the variability caused by combining noise and microelectrode data from different tissues. Nevertheless, the estimate is reasonable since the extrapolated reversal for \( i_{KBa} \) occurs near electrochemical equilibrium (i.e., \( \Delta \mu_K^i/F \approx 3 \text{ mV} \)).

### Density of Basolateral K Channels

With a simple on-off blocking scheme, the probability that a single channel is unblocked (\( P_o \)) is

\[
P_o = k_{10} / (k_{01}[Ba] + k_{10}) = k_{10} / 2\pi f_c.
\]

From Eqs. 6 and 14, the number of channels per square centimeter tissue area (i.e., channel density, or \( M \)) can be determined as

\[
M = I_{KBa} 2\pi f_c / i_{KBa} k_{10}.
\]

Channel density was calculated using the values for \( I_{KBa} \), \( i_{KBa} \), \( f_c \), and \( k_{10} \) from five recta at low [Ba] (0.5–1.0 mM), where \( k_{10} \) could be obtained from the \( 2\pi f_c [Ba] \) plot. We calculated a density of 178 ± 17 x 10⁶ channels/cm².
Locust rectal cells are columnar and the basolateral membrane is extensively infolded. According to electron micrographs (Phillips, J. E., M. S. Jarial, and H. B. Irvine, unpublished observations), the actual area of basolateral membrane is at least 20-fold greater than the macroscopic tissue area. Assuming 20-fold area amplification, we would expect a density of less than one channel per square micron of basolateral membrane, or <200 channels per cell. This estimate is somewhat higher than that deduced for Ca-activated K channels in pancreatic acinar cells (45 ± 11 per cell, obtained by comparing whole cell and single channel current; Maruyama et al., 1983b); however, a greater number of K channels would not be surprising considering the size of locust rectum cells and the importance of K absorption to insect renal function.

**DISCUSSION**

The locust rectum resembles most tight epithelia by having a significant basolateral membrane K conductance (Hanrahan and Phillips, 1984b). The results described in this paper indicate that the conductance results from Ba-sensitive channels. Using noise analysis and microelectrodes, we obtained evidence that the Ba association rate constant depends on membrane voltage, and the attenuation of transepithelially measured current fluctuations by apical membrane resistance was quantified. After making appropriate corrections for these effects, single channel current and ion-sensitive microelectrode data were used to estimate the conductance and density of basolateral K channels.

**Comparison of Single Channel and Basolateral Membrane Conductance**

Previous electrophysiological studies of locust rectum were not performed under the conditions used here (i.e., cAMP-stimulated, Cl-free, and high serosal [K]).
However, with the assumption that serosal Ba inhibits only basolateral membrane conductance, we can estimate individual membrane resistances under these conditions by plotting transepithelial conductance ($G_t$) vs. $\alpha/(1 + \alpha)$ at various Ba concentrations. Junctional and apical membrane conductances can be calculated from the y-intercept and slope, respectively, where

$$G_t = G_j + G_a\left[{\alpha}/(1 + \alpha)\right].$$

(16)

Linear regression of the data at 0, 1, 5, and 10 mM Ba (Figs. 2C and 8) gave a line with $r^2 = 0.983$, slope ($G_a$) = 3.47 mS cm$^{-2}$, and intercept ($G_j$) = 1.25 mS cm$^{-2}$, from which we calculated $R_b = 182 \Omega$ cm$^2$ in the absence of Ba. Not surprisingly, this value is lower than that measured previously by cable analysis when serosal [K] was only 10 mM (429 $\Omega$ cm$^2$; Hanrahan and Phillips, 1984b). $R_b$ increased from 163 to $>1,500 \Omega$ cm$^2$ when [Ba] was elevated to 10 mM. Paracellular resistance ($R_j = 1/G_j$) is 800 $\Omega$ cm$^2$, which is comparable to that obtained previously using flat-sheet cable analysis of cAMP-stimulated recta bathed with Cl-free saline (1,047 $\Omega$ cm$^2$; Hanrahan and Phillips, 1984b).

Calculating $G_a$ and $G_j$ in this way assumes that both are unaffected by the addition of Ba to the serosal side. This assumption seems reasonable because 20 mM Ba had no effect on $R_b$ when added to the mucosal side, nor did mucosal Ba affect the $I_o$ or power spectra. A large effect of intracellularly trapped Ba on the apical membrane seems unlikely in view of the dramatic (10-fold) reduction in $\alpha$ caused by adding Ba to the serosal side. Moreover, the apparent voltage dependence of Ba block (i.e., an increase in the on rate after hyperpolarization) would be consistent with block by external but not intracellular Ba.

Do microscopic estimates of single channel conductance and channel density agree with the macroscopic conductance of the basolateral membrane? With a density of 178 million channels/cm$^2$ and a single channel conductance of 60 pS (Fig. 8), we would predict the macroscopic K conductance to be $\sim$10.7 mS cm$^{-2}$. This value is higher than that calculated above ($\sim$5 mS cm$^{-2}$), but the discrepancy could be explained if the probability of channels being in the open state in the absence of Ba was $\sim$0.5. Noise resulting from rapid, spontaneous fluctuations may have been missed because we did not analyze frequencies above 100 Hz. Using the patch-clamp technique, K channels have been identified in the basolateral membrane of exocrine acinar cells (Maruyama et al., 1983a) and in an absorptive epithelium, rabbit urinary bladder (Hanrahan et al., 1984a). In both preparations, single channel conductance is $\sim$200 pS when patches are bathed symmetrically by $\sim$150 mM KCl, and spontaneous open-closed transitions are voltage dependent. In rabbit bladder, the open state probability decreases from 0.71 to 0.12 as the membrane is hyperpolarized from -20 to -80 mV. If insect basolateral K channels are similar, it is conceivable that they are open $\sim$50% of the time under the conditions used in this study.

**Potential Effects of Ba-insensitive Conductance in the Basolateral Membrane**

According to the circuit model in Fig. 1C, all of the Ba-insensitive $I_a$ measured under our experimental conditions must flow through the paracellular pathway. Consequently, the total current through Ba-sensitive channels at any particular [Ba] can be obtained by subtracting the $I_a$ that remains after adding a saturating
dose of Ba (i.e., $I_R$ in Eq. 4). However, if $I_R$ is also comprised of a Ba-insensitive current flowing through the basolateral membrane (e.g., a leak in parallel with the Ba-sensitive channels), then subtraction of $I_R$ would "over-correct" $I_K$ at low blocker concentrations, where the membrane potential, and hence the leak current, would be smaller. We estimated the largest errors in $I_{KBa}$ ($6 \mu A/cm^2$) and $i_{KBa}$ (0.14 pA) at 1 mM that could be expected from this effect, although it was necessary for the calculation to assume that Ba-sensitive conductance is completely blocked at 10 mM Ba, so that all the basolateral current is leak at this [Ba]. With this "worst-case" assumption, $g_{KBa}$ was calculated to be 48 pS using Eq. 13 (vs. 60 pS). Nonetheless, the effect of basolateral leak pathways on our estimate of single channel conductance is probably much less than 12 pS because there is still significant Ba-sensitive current remaining in the presence of 10 mM [Ba] (Figs. 2C and 3C).

**Why Does Adding Ba Hyperpolarize the Basolateral Membrane?**

Under short-circuit conditions, the apical and basolateral membranes are effectively in parallel, so that the intracellular potential is

$$V_b = E_a R_a/(R_a + R_b) + E_b R_b/(R_a + R_b).$$

(17)

K gradients determine $E_a$ and $E_b$ because both membranes are highly K selective under Cl-free conditions (Hanrahan and Phillips, 1984b). It can be seen from Table I that $E_b$ is much smaller (less negative) than $E_a$ because of the presence of high-K saline on the serosal side. According to Eq. 17, increases in $R_b$ caused by the addition of Ba will diminish the contribution of $E_b$ (approximately −5 mV), thereby driving $V_b$ toward $E_a$ (approximately −50 mV). Thus, simple changes in basolateral conductance can account for the hyperpolarization, although some contribution of other factors, such as the Na/K pump, cannot be excluded.

**Interpreting the Voltage Dependence of Ba Block**

A voltage-induced increase in the Ba association rate constant provides the most straightforward explanation for the curvature in the $2\pi f_c[Ba]$ plot, and is consistent with findings from other preparations. Nonetheless, three other possible mechanisms should also be mentioned. First, if more than one Ba ion binds to the K channel, the $2\pi f_c[Ba]$ plot might be described by a power function with a curvature resembling that in Fig. 6. However, it is not obvious why more than one Ba ion would be required to block the pore and other K channels have been shown to bind Ba in a one-to-one manner (e.g., Armstrong et al., 1982; Vergara and Latorre, 1983). A second possibility is that the binding of Ba to one channel might directly increase the affinity of a neighboring channel for Ba independently of membrane voltage. There is no precedent for this type of cooperative block, although it cannot be excluded at this time. A third possibility is that Ba actually passes through the channel so that the increase in $2\pi f_c$ with voltage could be due partly to enhanced entry into the cell (effectively an increase in the off rate). These alternatives, which require more assumptions than our hypothesis of voltage-dependent association, will be tested in the future using single channel methods.
The electrical distance $\gamma$ of 0.76 suggests that the barrier to externally applied Ba is deep within the membrane electric field. This implies an asymmetrical barrier and suggests that $k_{i0}$ should be less voltage dependent than $k_{o1}$. A similar asymmetry has already been shown for block by internal Ba of Ca-activated K channels incorporated into planar bilayers, where the association rate constant from the cis side (i.e., cytoplasmic side) depends on voltage ($\gamma = 0.65$) but the dissociation rate constant does not (Vergara and Latorre, 1983).

Under the conditions used in this study, voltage dependence should result in a lowering of the macroscopic $K_i$ for Ba at high concentrations; the affinity of Ba for the channel should appear to increase as [Ba] is elevated. To test this possibility, $K_i$ was calculated from the mean currents at each [Ba] shown in Fig. 2C, assuming that Ba inhibition is Michaelis-Menten at any particular Ba concentration. Fig. 10 shows that the macroscopic $K_i$'s generally declined as [Ba] was elevated. Moreover, when $K_i$ and $2\pi\sigma_f$ were used to solve for $k_{o1}$ (assuming $K_i = k_{i0}/k_{o1}$), $k_{o1}$ was found to increase with [Ba]. These calculations ignore possible competition between K and Ba (e.g., Sperelakis et al., 1967) but they provide independent evidence that the $2\pi\sigma_f$-[Ba] curvature can be explained by an increase in $k_{o1}$.

A deep binding site for extracellular Ba has been reported for the delayed rectifier in squid axon ($\gamma = 0.66$, Armstrong et al., 1982) and the inward rectifier in muscle (0.7, Standen and Stanfield, 1978). By contrast, external Ba shows
much weaker voltage dependence in the Ca-activated K channel from rabbit muscle (0.35, Ba on the trans side, Vergara and Latorre, 1983). An apparent distinction between these channels is the ability to contain more than one ion. Delayed and inward rectifiers are multi-ion pores (Hodgkin and Keynes, 1955; Hille and Schwarz, 1978), whereas Ca-activated K channels are thought to contain only a single ion (Vergara and Latorre, 1983; but see Latorre et al., 1986). The strong voltage dependence of block by external Ba would seem more consistent with basolateral K channels being multi-ion pores, and we note that Kirk and Dawson (1983) reached the same conclusion from tracer studies of turtle colon basolateral membrane. They observed interactions between “K-like” cations and deviations from Ussing’s flux ratio (Ussing, 1949) when transepithelial ⁴²K fluxes were measured after permeabilizing the apical membrane and inhibiting the Na pump. A flux ratio exponent of ~2 was obtained, which for a single-file model is related to the number of ions within the channel (Hodgkin and Keynes, 1955; see Heckmann, 1972). Armstrong et al. (1982) have pointed out that the electrical distance for the binding of voltage-dependent blockers such as Ba is probably not unique in multi-ion pores and might be voltage dependent. It is therefore conceivable that the blocking site for Ba may be slightly deeper when the membrane is at the normal potential of ~50 mV.

**Ba Sensitivity**

Since the early demonstration by Pacifico et al. (1969) that serosal Ba inhibits basolateral K conductance in gastric mucosa, evidence for Ba-sensitive permeability has been obtained in at least nine epithelia and the list continues to grow (cf. Lewis et al., 1984b). Sensitivity to serosal Ba appears to be highest in frog skin (Nagel, 1979), intermediate in trachea (Welsh, 1983), and lowest in locust rectum. The concentration required for half-maximal inhibition (\( K_\text{b} \)) varies between 0.2 and 2.0 mM in these tissues, but this is a reasonable range considering the experimental conditions under which Ba was tested. Different membrane potentials might partly explain the variation, as would different external [K] concentrations: basolateral membrane potential was highest in the frog skin experiments (ranging from ~95 to ~45 mV after Ba addition), intermediate in trachea (~51 to ~28 mV), and lowest in the present study (~8 to ~26 mV). We would expect K-Ba competition to have been lowest for frog skin (2.5 mM K), intermediate for trachea (5.4 mM K), and highest in the present study (80 mM K; high-K saline). In comparing Ba inhibition in epithelial and excitable cells, it is interesting that half-maximal block in the millimolar range is similar to findings for delayed rectifiers in frog muscle (0.8 mM Ba; Standen and Stanfield, 1978) and squid axon (3 mM Ba at ~70 mV; Armstrong et al., 1982), and also for Ca-activated K channels from rabbit muscle (1.8 mM Ba on trans side at 0 mV; Vergara and Latorre, 1983).

**Selectivity and Conductance**

The ratio \( P_{\text{Na}}/P_{\text{K}} \) cannot be estimated reliably from the noise data presented here; however, there is evidence that basolateral K channels are highly selective against Na. First, replacing serosal Na by K (under Cl-free conditions, with K activity elevated approximately sevenfold) caused \( V_i \) to hyperpolarize by ~50 mV
(Fig. 2A). Second, the net electrochemical driving force for K is normally <4 mV above equilibrium despite high serosal [Na] (i.e., >100 mM; Hanrahan and Phillips, 1984a,b). Selectivity sequences reported for basolateral K channels generally resemble those of excitable cells: Tl and Rb can both penetrate the basolateral K pathway in turtle colon ($P_{TL} \approx P_K$, $P_{Rb} = 0.1 P_K$; Kirk and Dawson, 1983). Also, Cs inhibits inward K current at the basolateral membrane of both rabbit colon (Wills et al., 1979) and locust rectum (Hanrahan, J. W., and J. Meredith, unpublished observations).

Locust hindgut K channels appear to have intermediate conductance, lower than many Ca-activated K channels (~180–220 pS in ≥100 mM KCl) and the sarcoplasmic reticulum K channel (~120 pS), but higher than delayed and inward rectifiers (2–18 pS, various conditions; see review by Latorre and Miller, 1983). This preliminary estimate of 60 pS for single channel conductance needs to be confirmed eventually using the patch-clamp technique. It is likely that basolateral K channels would have a much larger conductance than 60 pS if they were bathed symmetrically with higher K concentrations as in studies of isolated channels. For comparison, when K channels from sarcoplasmic reticulum are incorporated into lipid bilayers, their conductance is half-maximal in solutions containing 50 mM K activity (Coronado et al., 1980). A similar conductance-activity relationship for locust K channels would result in single channel conductances of ~120 pS when they are bathed symmetrically with 150 mM [K]. Finally, we suspect that locust K channels are Ca insensitive, like those from rabbit urinary bladder, because maneuvers that should drastically alter intracellular [Ca] have no effect on transepithelial conductance (Hanrahan and Phillips, 1985).

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